Periodic homogenization and harmonic measures

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Introduction: L^p -solvability of the Dirichlet problem Consider the Dirichlet problem

$$\begin{cases} -\nabla \cdot A \nabla u = 0 & \text{ in } \Omega, \\ u = f & \text{ on } \partial \Omega, \end{cases}$$

with A a uniformly elliptic and bounded matrix field and Ω the upper-half plane, that is, $\Omega = \mathbb{R}^{d+1}_+ := \{(t, x) \in \mathbb{R}^{d+1} : t > 0\}.$

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$$\begin{cases} -\nabla \cdot A \nabla v = \nabla \cdot A \nabla F & \text{ in } \Omega, \\ v = 0 & \text{ on } \partial \Omega, \end{cases}$$

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with F an H^1 -extension of f in Ω (in the trace sense). In this way, u = F + v is the solution to the original problem. This requires $f \in H^{\frac{1}{2}}(\partial \Omega)$ since Tr $H^1(\Omega) = H^{\frac{1}{2}}(\partial \Omega)$.

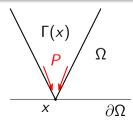
Introduction: L^p-solvability of the Dirichlet problem

Theorem ([Caf+81])

If u satisfies $Lu = -\nabla \cdot A\nabla u = 0$ in Ω , then the non-tangential limit

 $\lim_{\substack{P\in\Gamma(x),\\P\to x}}u(P)$

exists for almost every $x \in \partial \Omega$ wrt to the harmonic measure ω associated with the operator L and Ω .



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The L^p -solvability for some p is equivalent to $\omega \in A_{\infty}(\sigma)$, that is, the existence of $\varepsilon, \delta \in (0, 1)$ such that for any $E \subset Q \subset \partial \Omega$

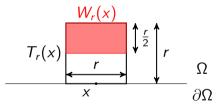
$$\frac{\sigma(E)}{\sigma(Q)} \leqslant \delta \quad \Rightarrow \quad \frac{\omega(E)}{\omega(Q)} \leqslant \varepsilon.$$

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When A = Id, that is, $L = -\Delta$, this property is verified. **Can we perturb the Laplacian and retain this property?** Define

$$T_r(x) := (0, r] \times Q_r(x), \quad W_r(x) := (r/2, r] \times Q_r(x).$$

Here $Q_r(x) \subset \partial \Omega = \mathbb{R}^d$ is the cube centered at x of side-length r.



Define the (local) oscillation of A at point $(t, y) \in \Omega$ by

$$\Big(\operatorname{osc} A\Big)(t,y) := \sup_{Y,Y' \in W_t(y)} |A(Y) - A(Y')|.$$

Theorem ([KP01])

If there is a constant C>0 such that for all r>0 and $x\in\partial\Omega$ one has

$$\int_{\mathcal{T}_r(x)} |\operatorname{osc} A|^2(t,y) \frac{dydt}{t} \leqslant Cr^d,$$

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This implies that $\operatorname{osc} A$ vanishes at a certain speed when $t \to 0$:

$$\int_{T_r(x)} |\operatorname{osc} A|^2(t, y) \frac{dydt}{t} = \int_0^r \left(\int_{Q_r(x)} |\operatorname{osc} A|^2(t, y) dy \right) \frac{dt}{t}.$$

Therefore, the DKP condition can be viewed as a reasonable perturbation of the Laplacian.

Moreover, this result is sharp in the sense of convergence speed: If

$$\int_{\mathcal{T}_r(x)} a^2(t,y) \frac{dydt}{t} = +\infty,$$

then there exists a uniformly elliptic matrix field A such that $\operatorname{osc} A \sim a$ and the harmonic measure associated to the operator $-\nabla \cdot A \nabla$ does not admit the A_{∞} property.

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Question: Can homogenization be used to replace the assumption of vanishing oscillations in the DKP condition? **Answer**: Yes!

Main theorem

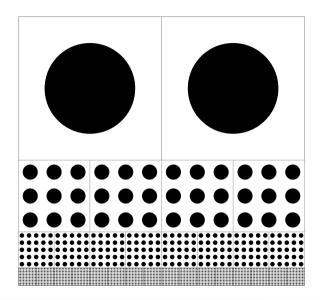
Suppose that $A_{\#}$ is a \mathbb{Z}^{d+1} periodic elliptic matrix field on \mathbb{R}^{d+1} with associated homogenized matrix $\bar{A}_{\#} = Id$. We define the coefficient field A in the following way

$$A(t,x) = egin{cases} \mathsf{Id} & t \geqslant 1, \ A_{\#}(rac{t,x}{2^k arepsilon_k}) & t \in [2^{k-1},2^k), k \leqslant 0. \end{cases}$$

Theorem ([Dav+25]) When $\varepsilon_k \lesssim 2^k$, $\omega \in A_{\infty}(d\sigma)$.

In particular, if $A_{\#}$ is not constant, the DKP condition cannot be satisfied.

Main theorem



A criterion for the A_∞ property of the harmonic measure is the following:

Theorem ([Ken+00])

The harmonic measure $\omega \in A_{\infty}(d\sigma)$ if and only if for any continuous solution u to the Dirichlet problem,

$$r^{-d}\int_{\mathcal{T}_r(x)}t|\nabla u(t,y)|^2dydt\leqslant C\|f\|_{L^{\infty}}^2.$$

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Ansatz:

$$u \sim u^{2s} := \bar{u} + \sum_{k \leq 0} 2^k \varepsilon_k \chi_k \phi_i(\frac{\cdot}{2^k \varepsilon_k}) \partial_i \bar{u},$$

where χ_k is a cut-off because of the boundary layer, \bar{u} is the solution to the Dirichlet problem with the same boundary datum f and $A = \bar{A}_{\#} = Id$.

The key estimate is that

$$r^{-d}\int_{\mathcal{T}_r(x)}t|\nabla z(t,y)|^2dydt\leqslant C\|f\|_{L^{\infty}}^2,$$

where $z := u - u^{2s}$ is the homogenization error. It satisfies the following equation:

$$-\nabla\cdot A\nabla z = \nabla\cdot f,$$

with

$$f := \sum_{k \leqslant 0} 2^k \varepsilon_k (A_\# \phi^i - \sigma^i)(\frac{\cdot}{2^k \varepsilon_k}) \nabla(\chi_k \partial_i \bar{u}) + \mathbf{1}_{t \in [2^{k+1}, 2^k]} (1 - \chi_k) (A_\# - \mathsf{Id})(\frac{\cdot}{2^k \varepsilon_k}) \nabla \bar{u}.$$

We localize the error in $T_r(x)$ by Caccioppoli's inequality, maximum principle and carefully choosing the boundary layer size:

$$\begin{split} \int_{\mathcal{T}_r(x)} |\nabla z|^2 &\lesssim \int_{\mathcal{T}_{2r}(x)} f^2 + r^{-2} \int_{\mathcal{T}_{2r}(x)} z^2 \\ &\lesssim \sum_{k \leqslant \min([\log_2 r], 0)} \varepsilon_k \int_{[2^{k+1}, 2^k) \times Q_r(x)} |\nabla \bar{u}|^2 + r^{d-1}. \end{split}$$

Therefore, if $\varepsilon_k \lesssim 2^k$, then one can deduce, by harmonicity of \bar{u} ,

$$\int_{\mathcal{T}_r(x)} \min(t,1) |\nabla z|^2 \lesssim \min(r,1) (\underbrace{\int_{\mathcal{T}_r(x)} t |\nabla \bar{u}|^2}_{\lesssim r^d} + r^{d-1}) \sim r^d.$$

- It's interesting to understand how the speed $\varepsilon_k \lesssim 2^k$ in the assumption can be interpreted or improved.
 - If we pick the coefficient field A differently in each layer [2^{k+1}, 2^k) (instead of a single A_#), then we need a stronger assumption ε_k ≤ 2^{3/2/k}.
 - If there exists a global solution for the corrector equation, then the weaker assumption $\varepsilon_k ≤ 2^{\frac{1}{2}k}$ is enough.
 - So Can we go further and prove the theorem for $\sum_{k \leq 0} \varepsilon_k^q < +\infty$?

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