

Modélisation de micro-filaments et convergence d'un modèle à N segments

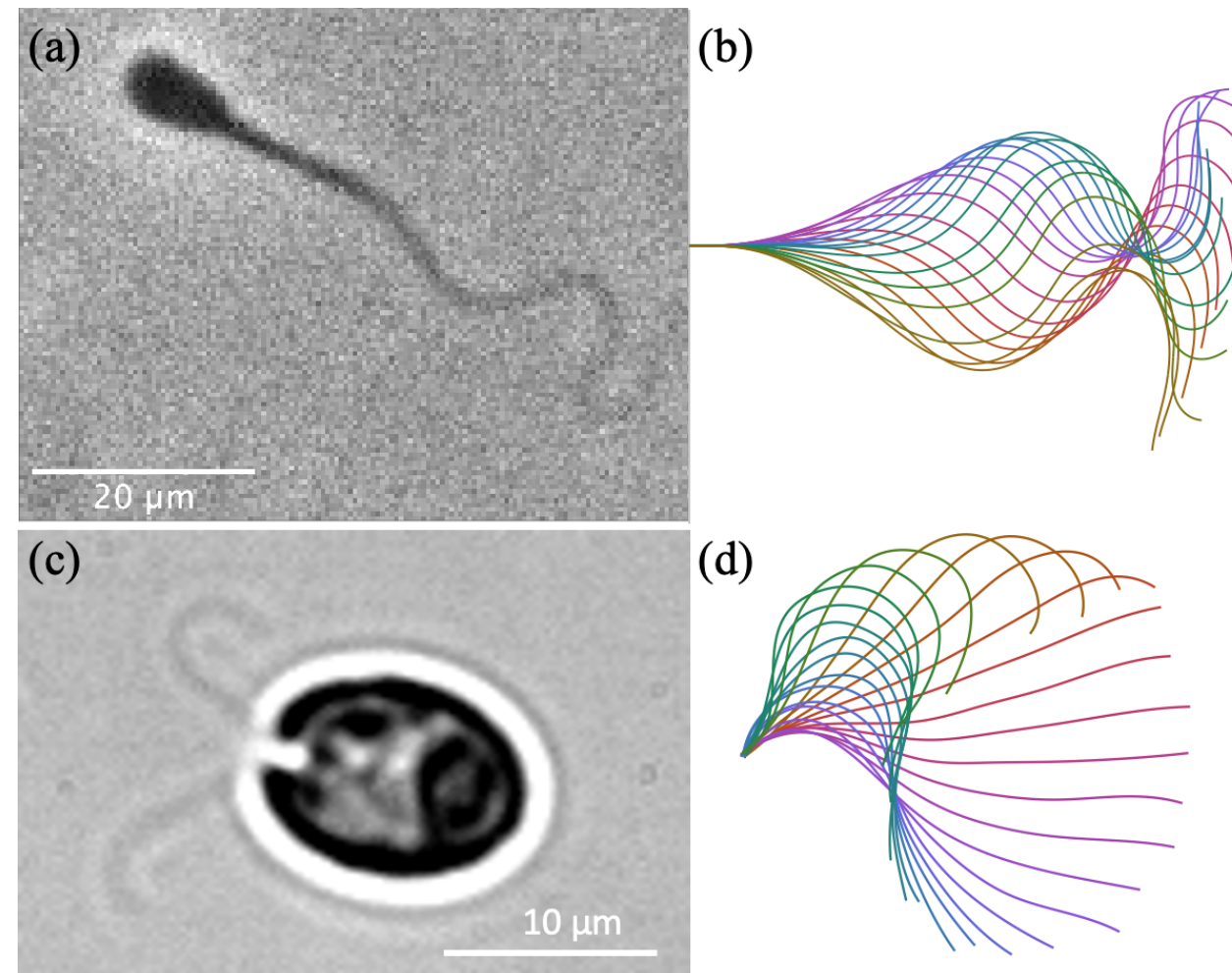
Clément Moreau



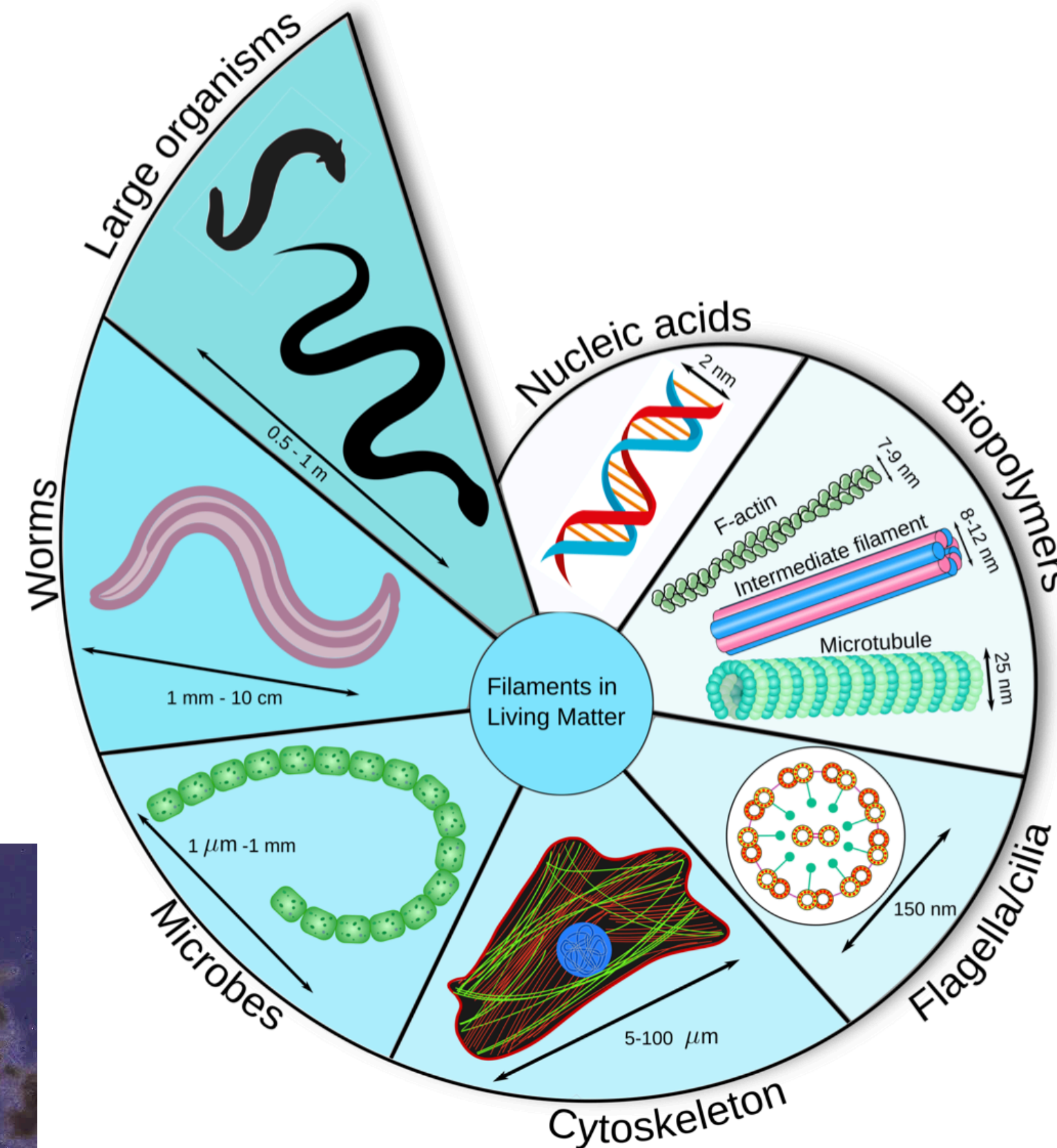
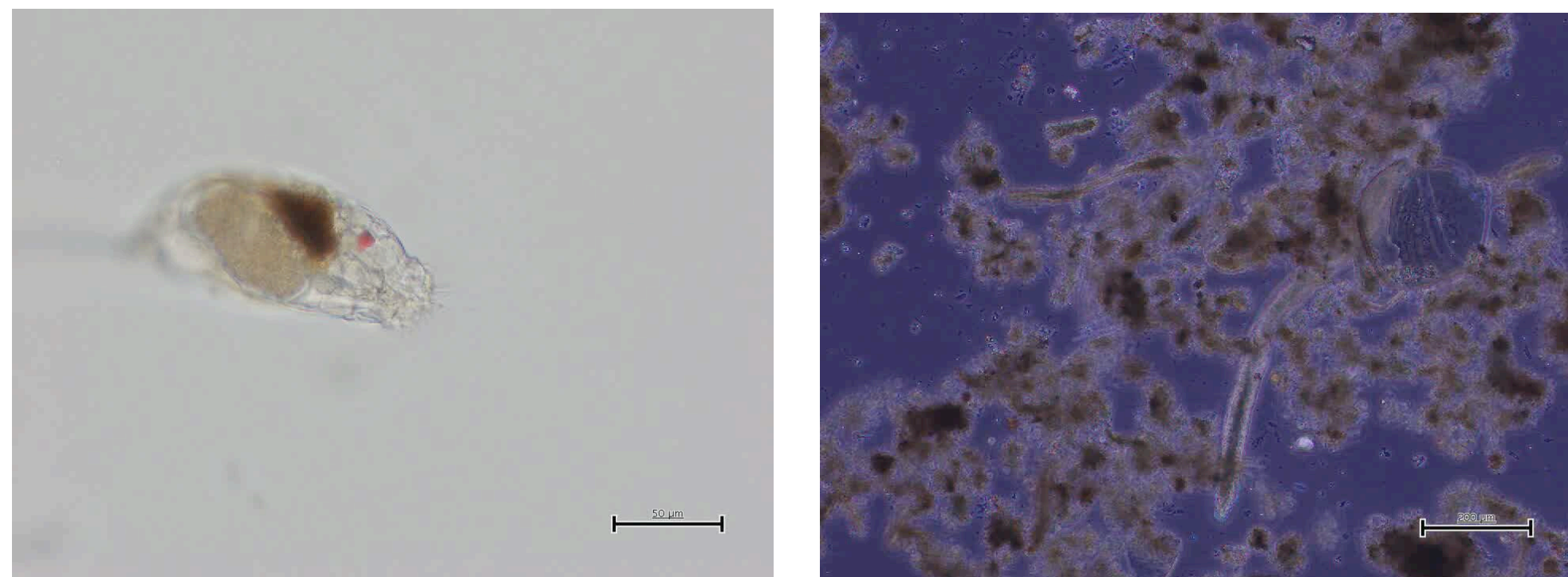
Congrès SMAI — 6 juin 2025

Flexible filaments are ubiquitous at microscopic scale

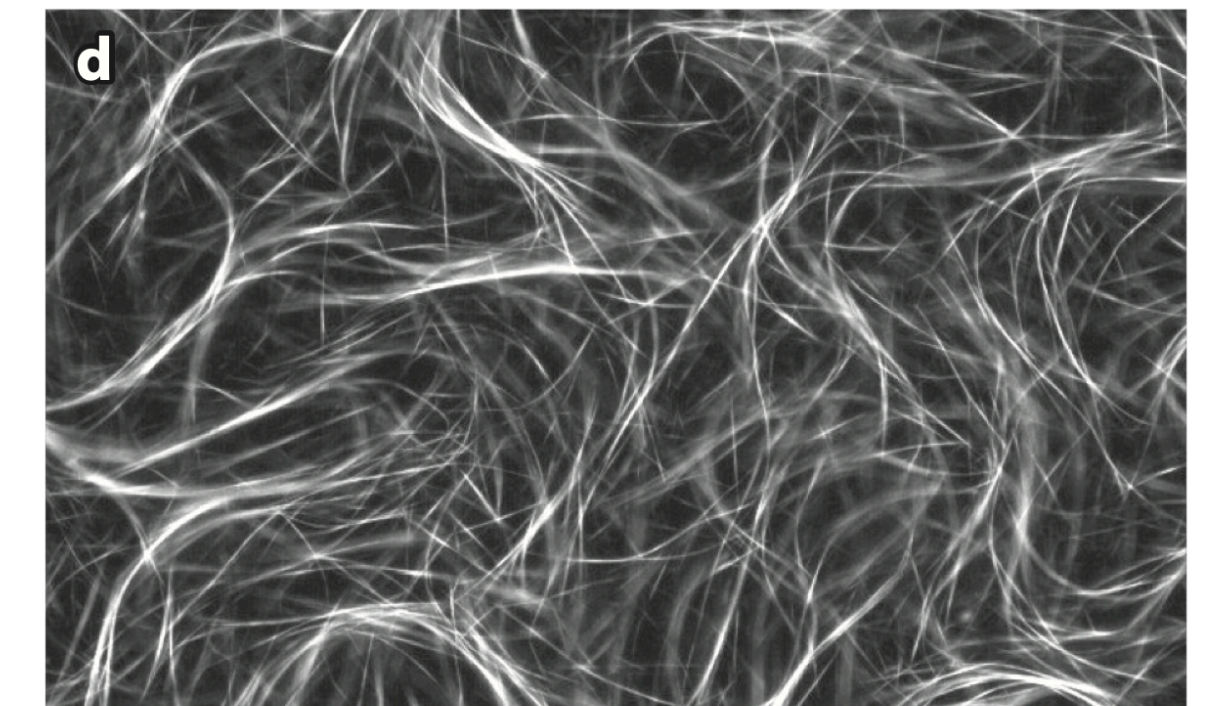
Moreau et al., Phys. Rev. Research 6 (2024)



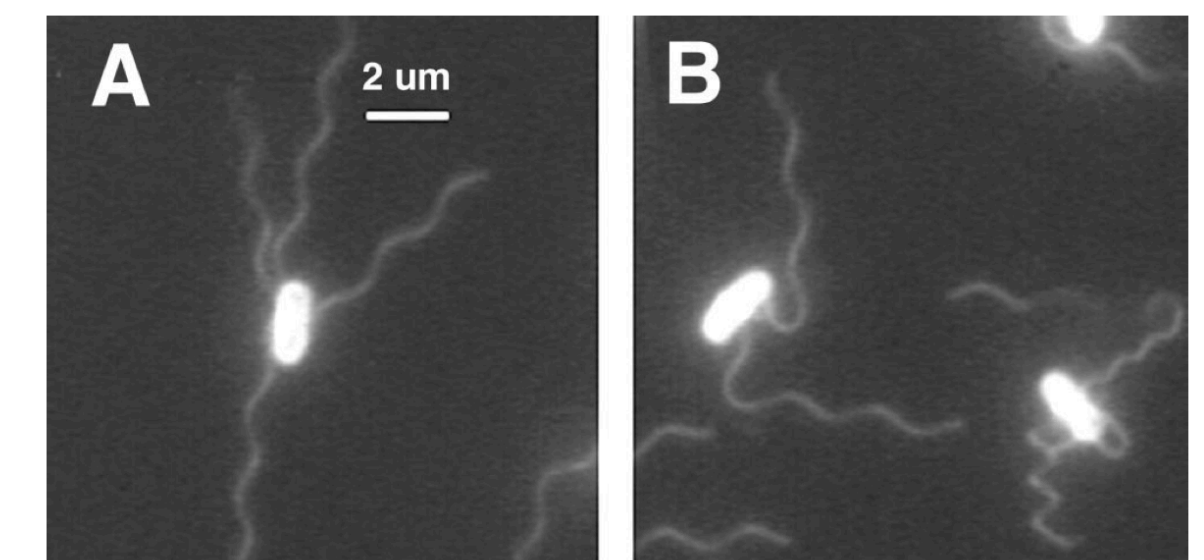
Kyoto University pond



Camman et al. arXiv:2502.12731



Sanchez et al., Nature 491 (2012)

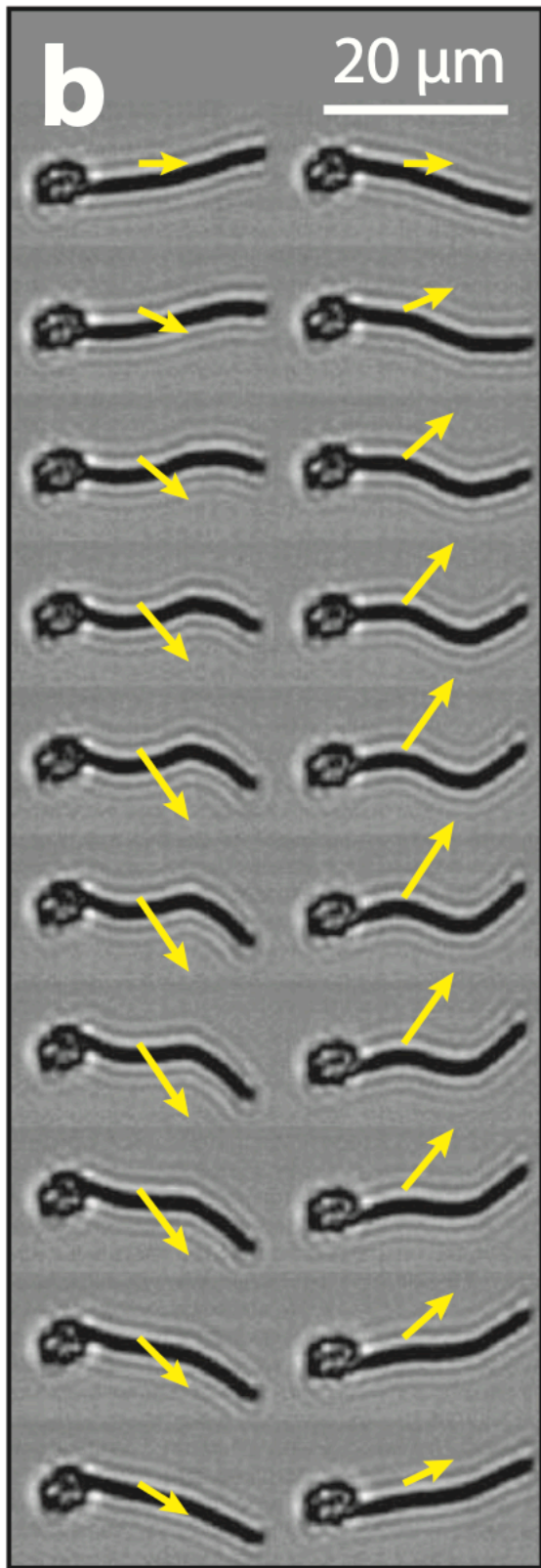


Turner et al., J. Bacteriol. 182 (2000)

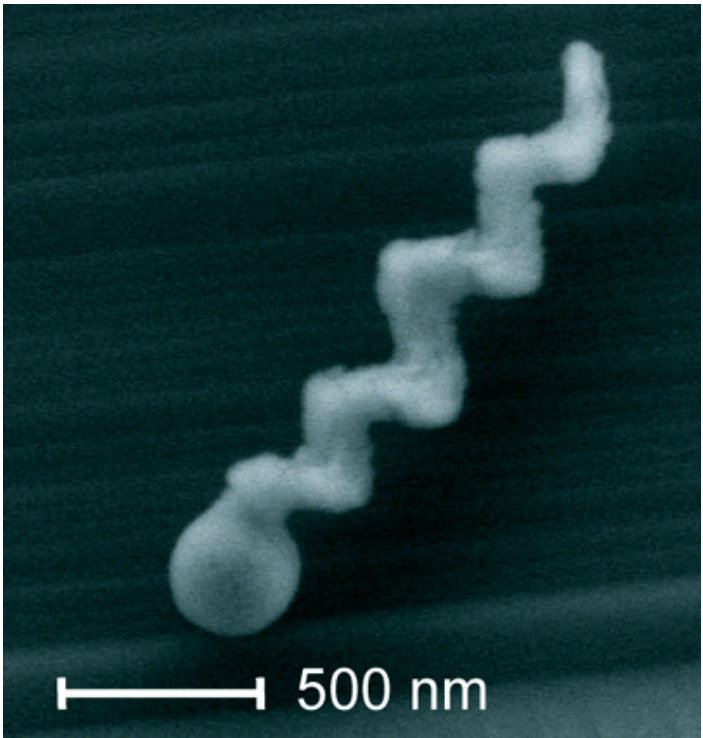
Dynamics : **elastic body** interacting with **fluid flow** (+ some kind of **activity**)

Bio-inspired slender micro-robots

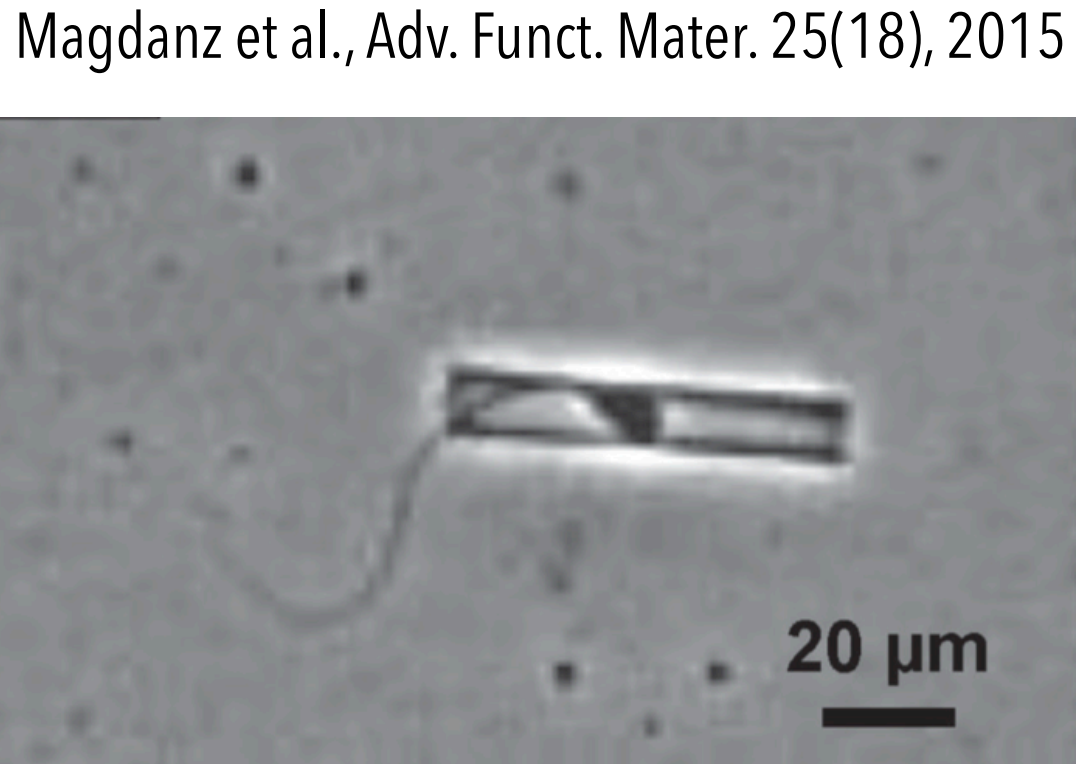
Palagi, Fischer, Nat. Rev. Materials 3 (2018)



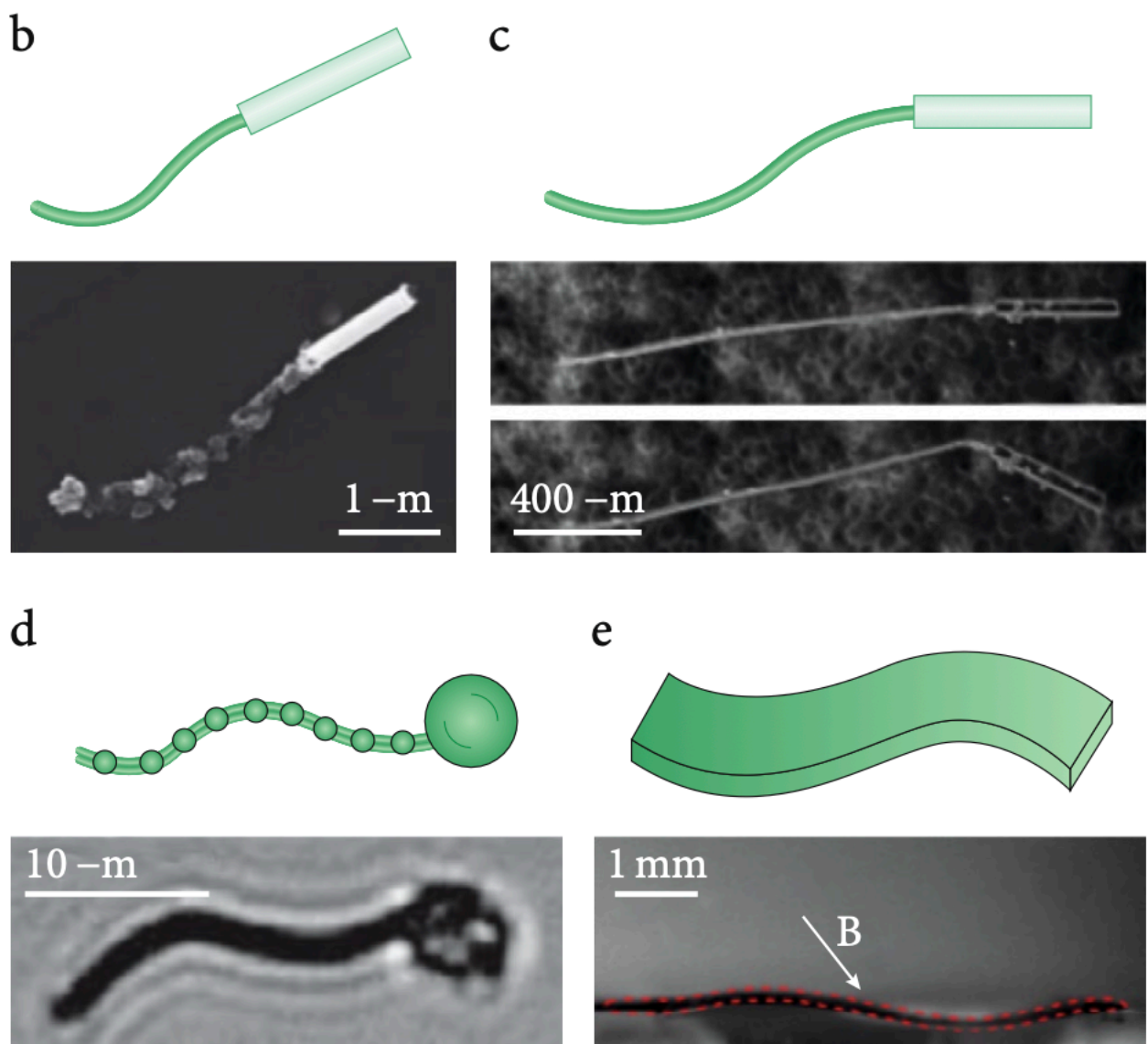
Xin et al., ACS Nano 15(11), 2021



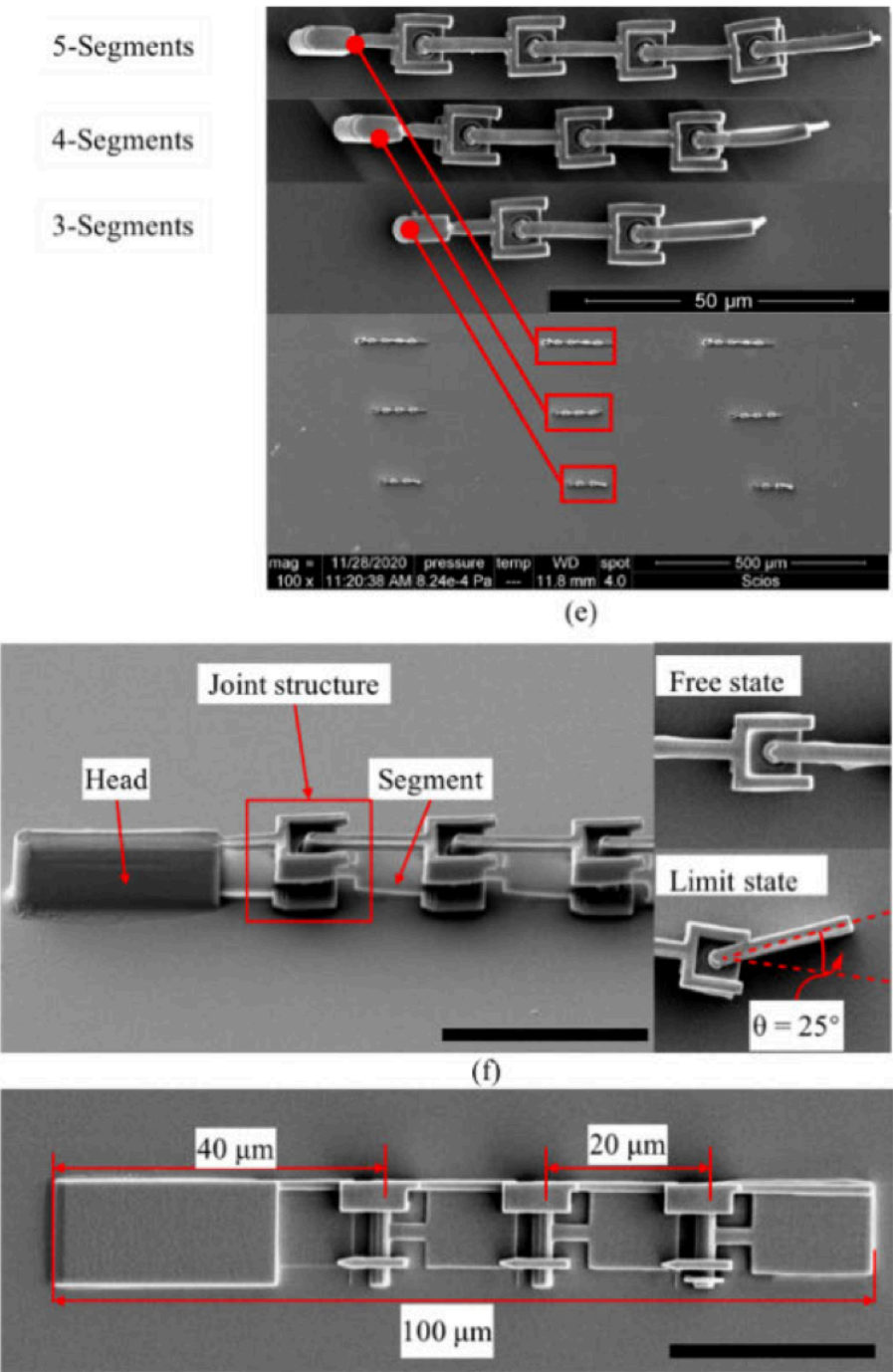
Ghosh & Fischer, Nano Letters 9(6), 2009



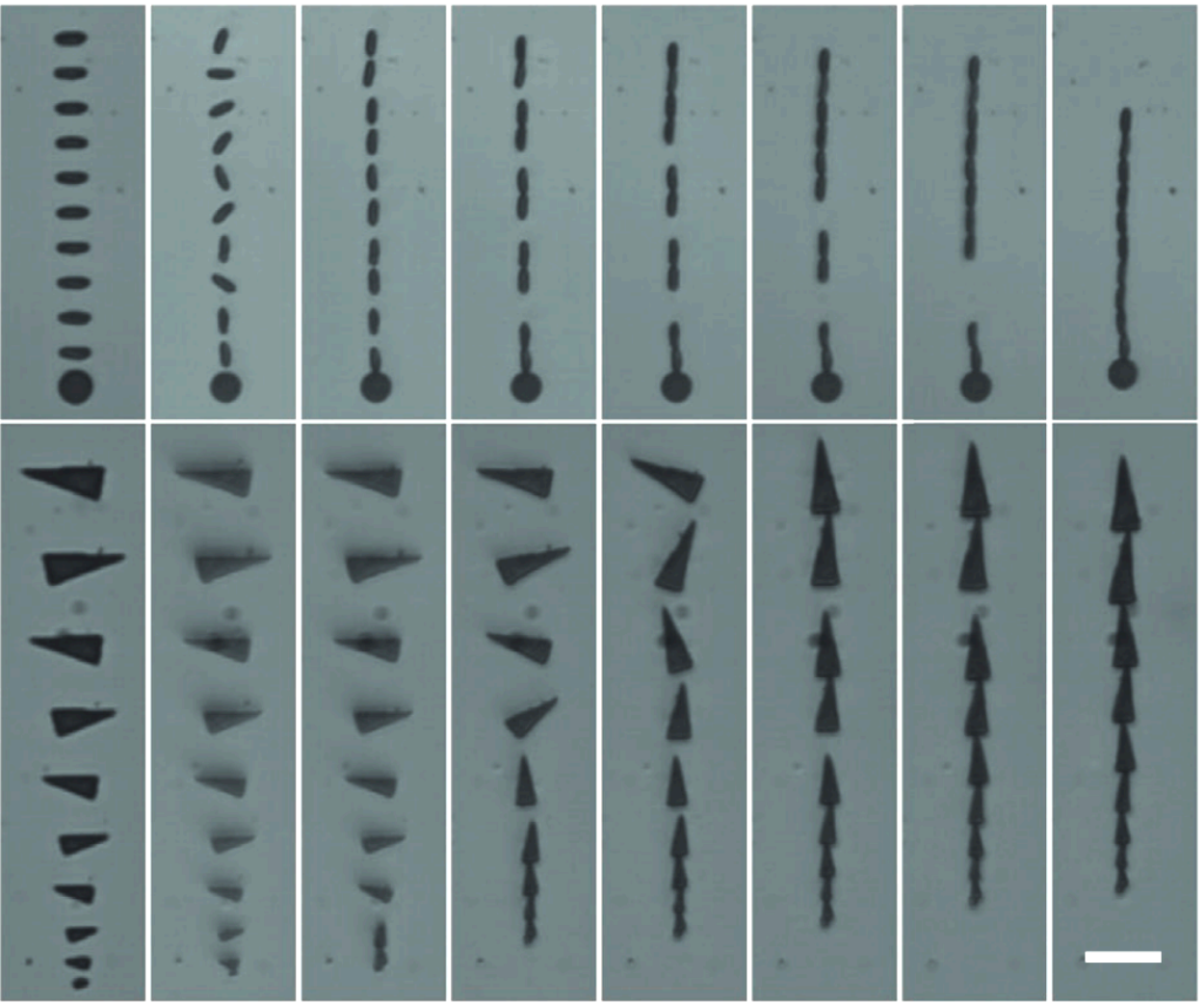
Magdanz et al., Adv. Funct. Mater. 25(18), 2015



Xing et al. IEEE Access 9, 2021



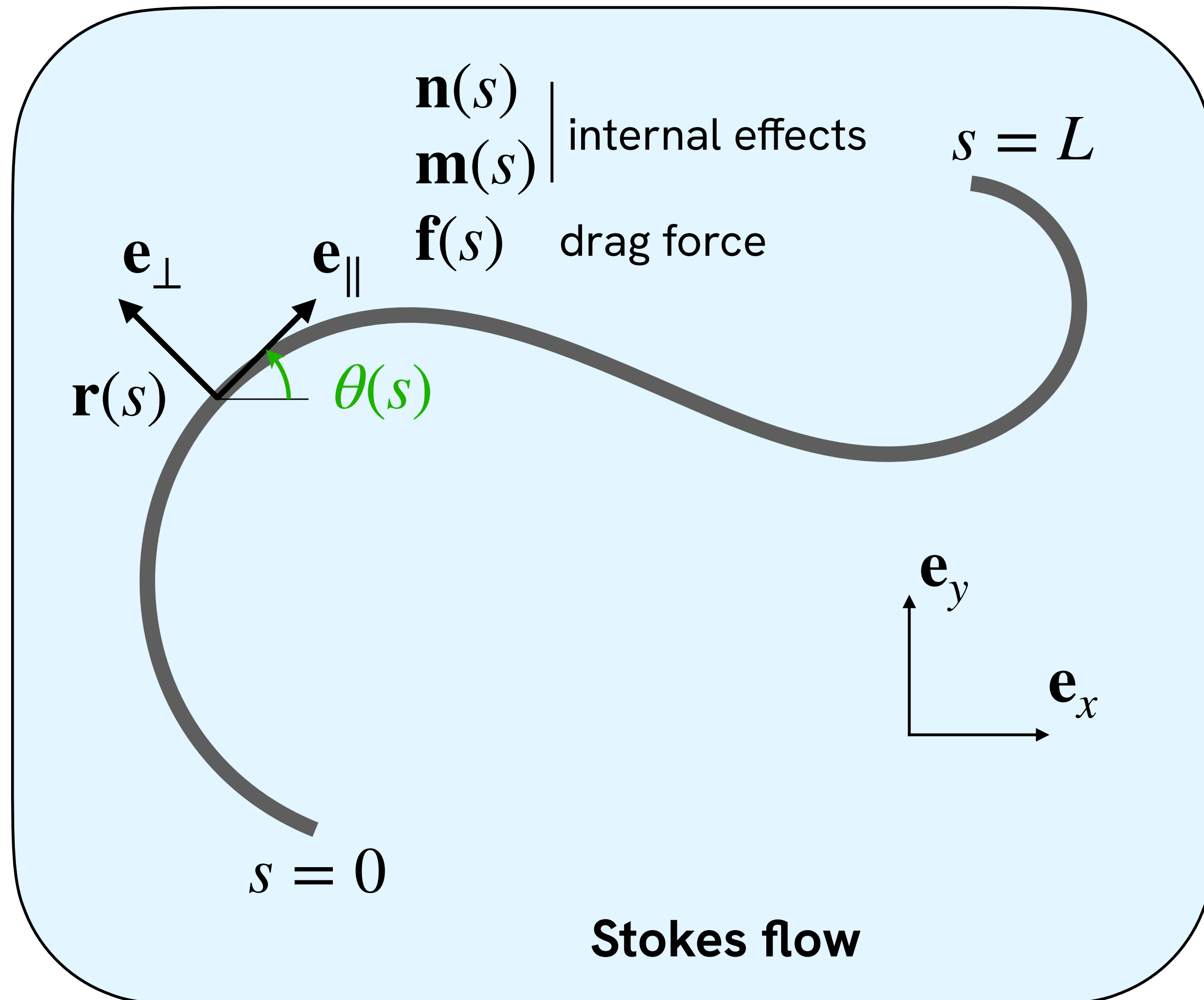
Pauer et al., Adv. Mater. 33, 2021



→ modelling and simulation

Elastohydrodynamics equations

Deformable slender body in a viscous fluid (Stokes flow)



Resistive Force Theory (RFT): local anisotropic operator

$$\mathbf{f}(t, s) = \mathbf{C}(\theta(t, s))\dot{\mathbf{r}} = -c_{\parallel}(\mathbf{e}_{\parallel} \cdot \dot{\mathbf{r}})\mathbf{e}_{\parallel} - c_{\perp}(\mathbf{e}_{\perp} \cdot \dot{\mathbf{r}})\mathbf{e}_{\perp}$$

hydrodynamic
force density

tangential
velocity

orthogonal
velocity

(Cont)

$$\partial_s \mathbf{n} + \mathbf{f} = 0 \quad (\text{local force balance})$$

$$\partial_s \mathbf{m} + \mathbf{r}_s \times \mathbf{n} = 0 \quad (\text{local torque balance})$$

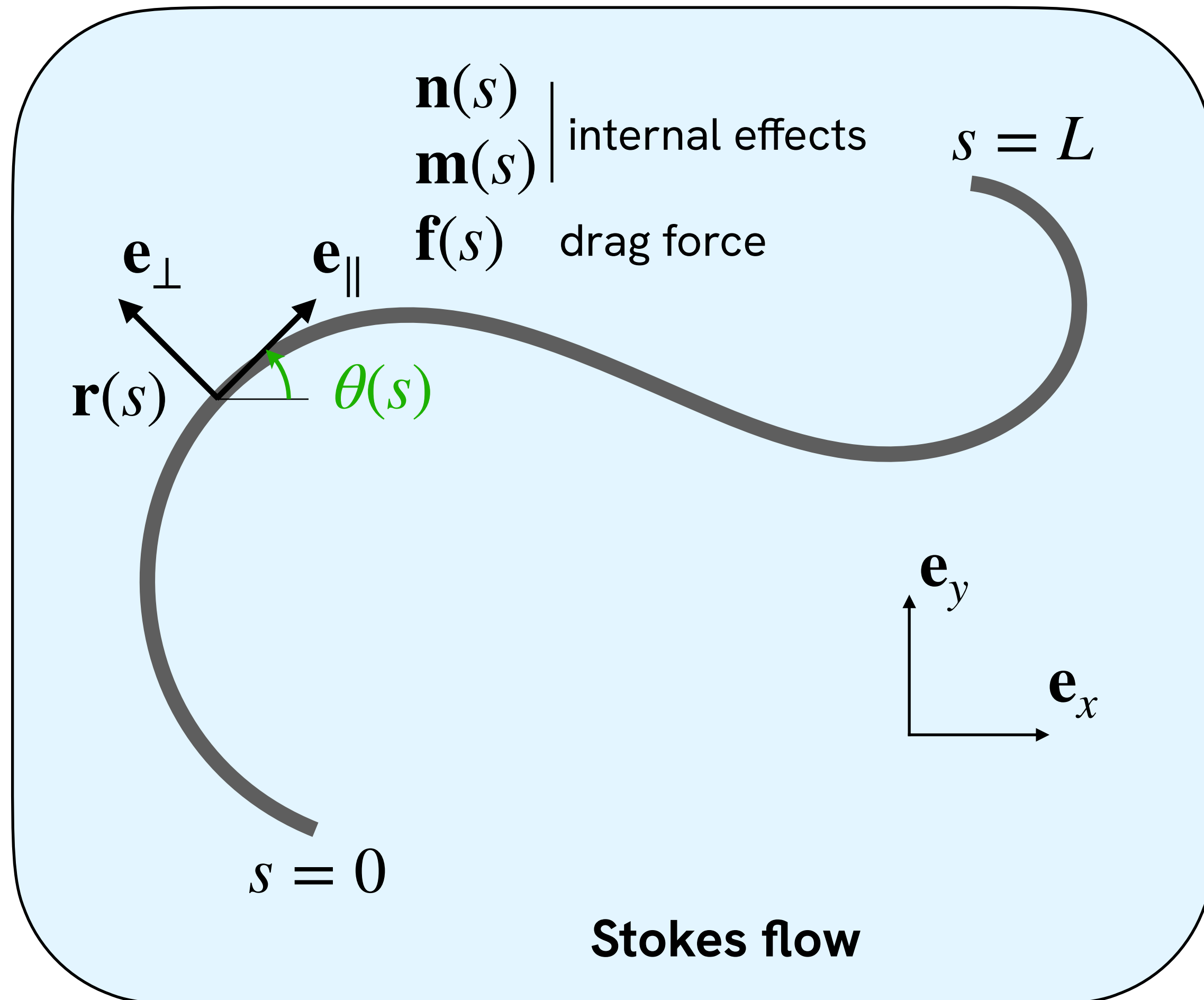
$$\mathbf{m} = E\theta_s \mathbf{e}_z \quad (\text{elastic constitutive law})$$

$$|\mathbf{r}_s|^2 = 1 \quad (\text{inextensibility})$$

+ B.C.

Elastohydrodynamics equations

Deformable slender body in a viscous fluid (Stokes flow)



(Cont)

$$\partial_s \mathbf{n} + \mathbf{C}(\theta) \partial_t \mathbf{r} = 0 \quad (\text{local force balance})$$

$$\partial_s \mathbf{m} + \mathbf{r}_s \times \mathbf{n} = 0 \quad (\text{local torque balance})$$

$$\mathbf{m} = E \theta_s \mathbf{e}_z \quad (\text{elastic constitutive law})$$

$$|\mathbf{r}_s|^2 = 1 \quad (\text{inextensibility})$$

+ B.C.

free

$$\mathbf{m} = 0, \mathbf{n} = 0$$

pinned

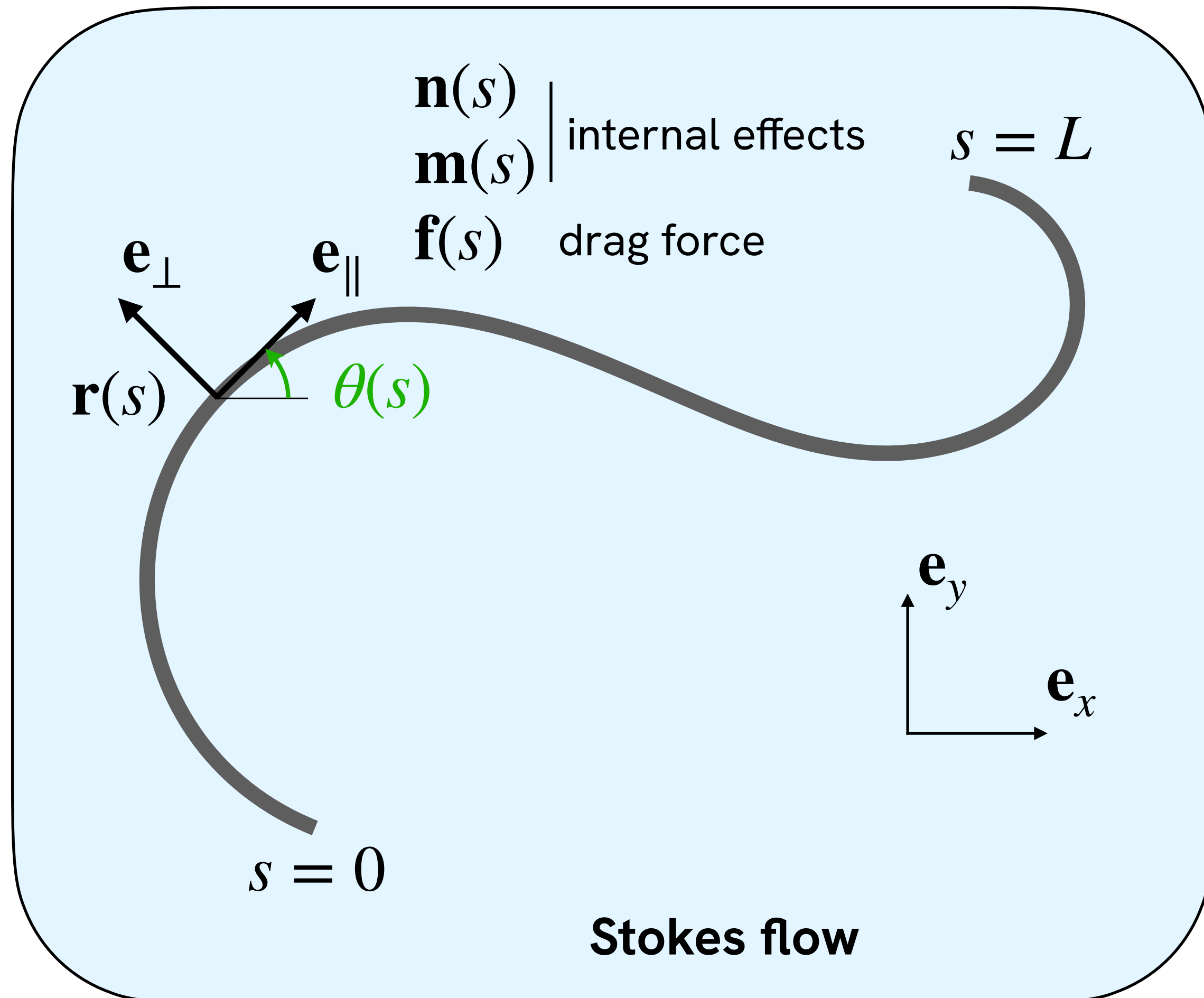
$$\mathbf{m} = 0, \dot{\mathbf{r}} = 0$$

clamped

$$\dot{\theta} = 0, \dot{\mathbf{r}} = 0$$

Elastohydrodynamics equations

Deformable slender body in a viscous fluid (Stokes flow)



(Cont)

$$\partial_s \mathbf{n} + \mathbf{C}(\theta) \partial_t \mathbf{r} = 0 \quad (\text{local force balance})$$

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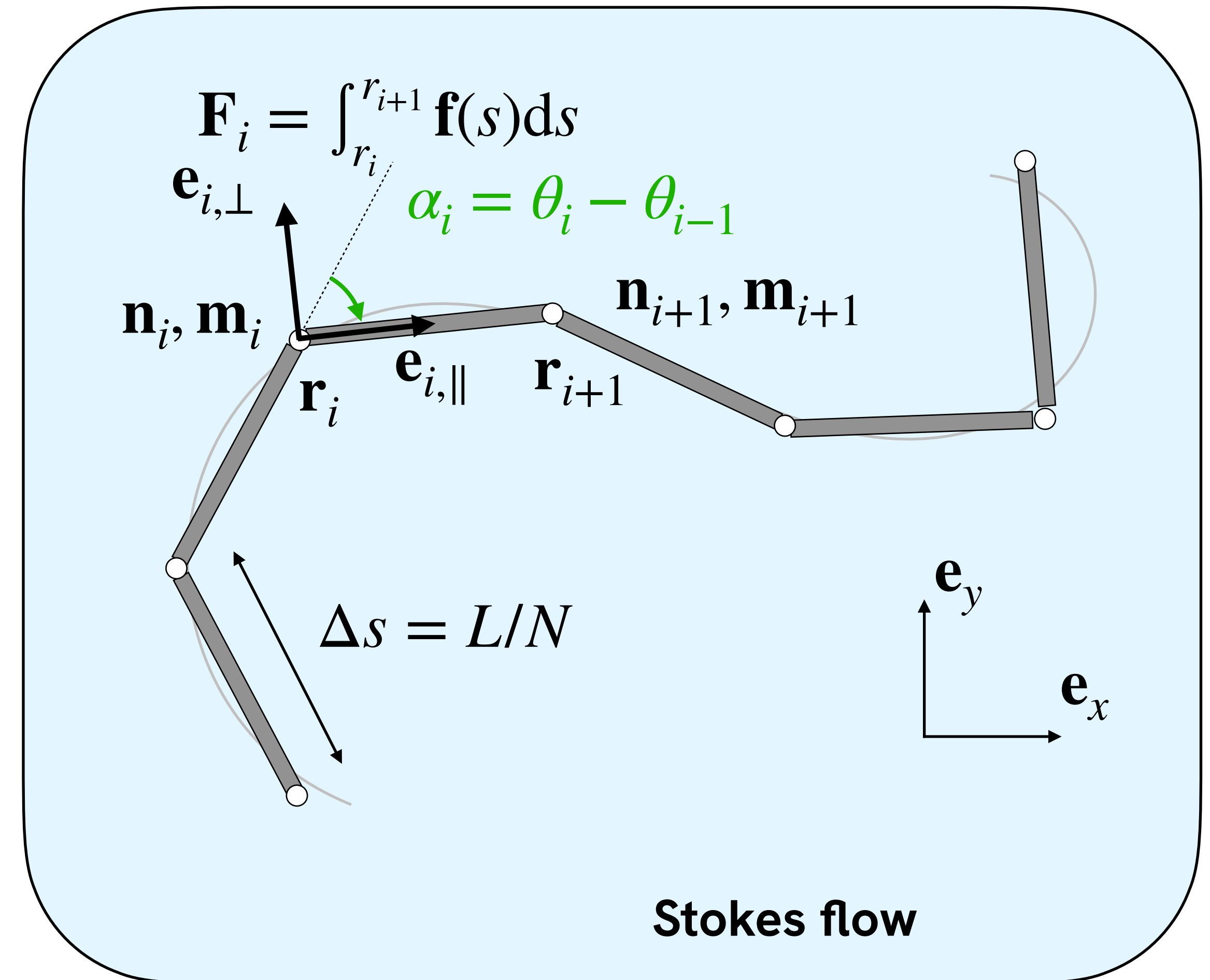
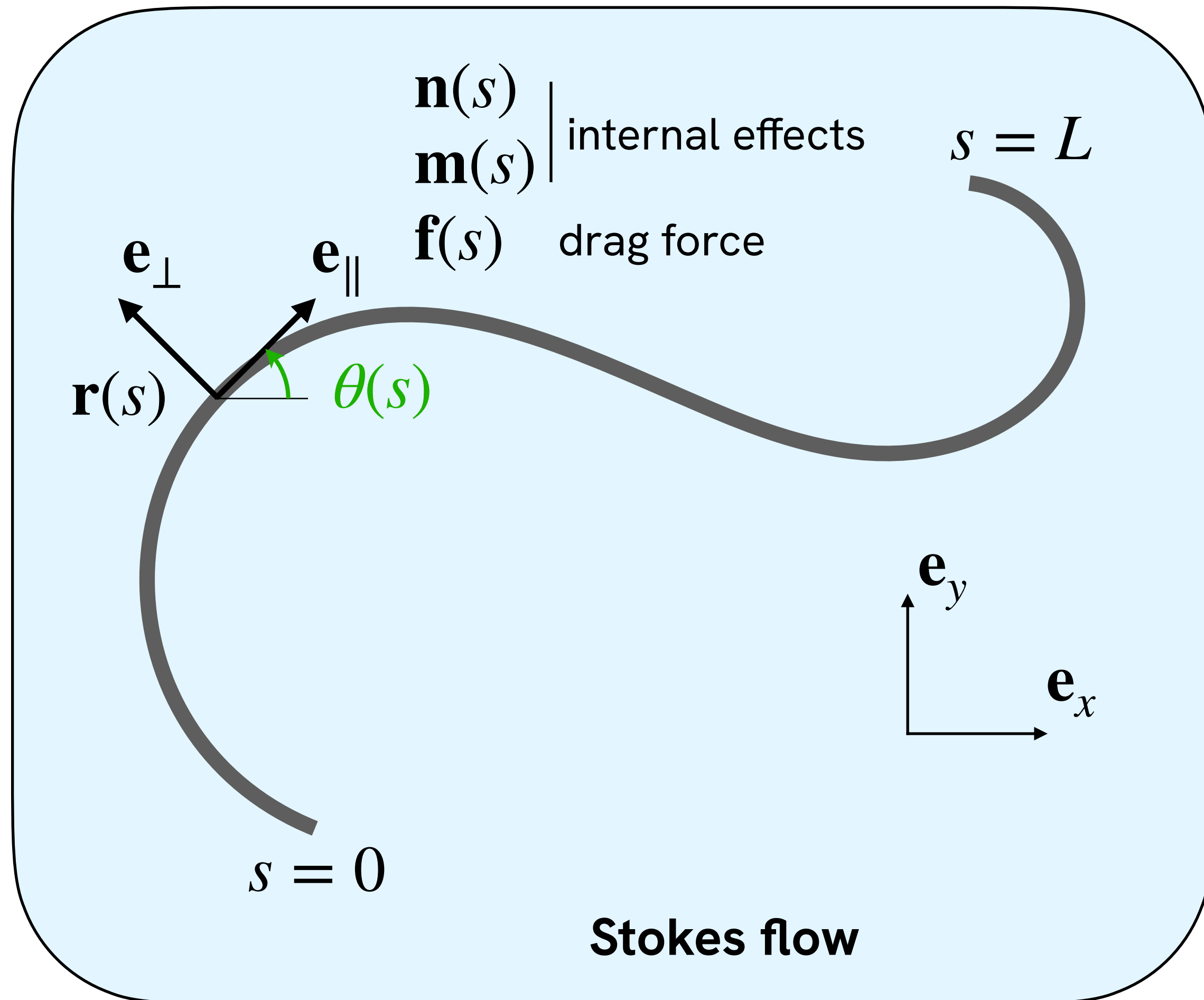
- **Nonlinear PDE** with 4th space derivative + τ (Lagrange multiplier for the tension)

$$\dot{\mathbf{r}} = -(\mathbf{I} + \gamma \mathbf{r}_s \mathbf{r}_s^T)^{-1} \cdot \partial_s (\mathbf{r}_{sss} - \tau \mathbf{r}_s)$$

- Challenging to establish **well-posedness** (existence and uniqueness of solutions) (Mori & Ohm, Nonlinearity 2023)
- Challenging to solve numerically (du Roure et al., Annu. Rev. Fluid Mech. 2017)

Elastohydrodynamics equations

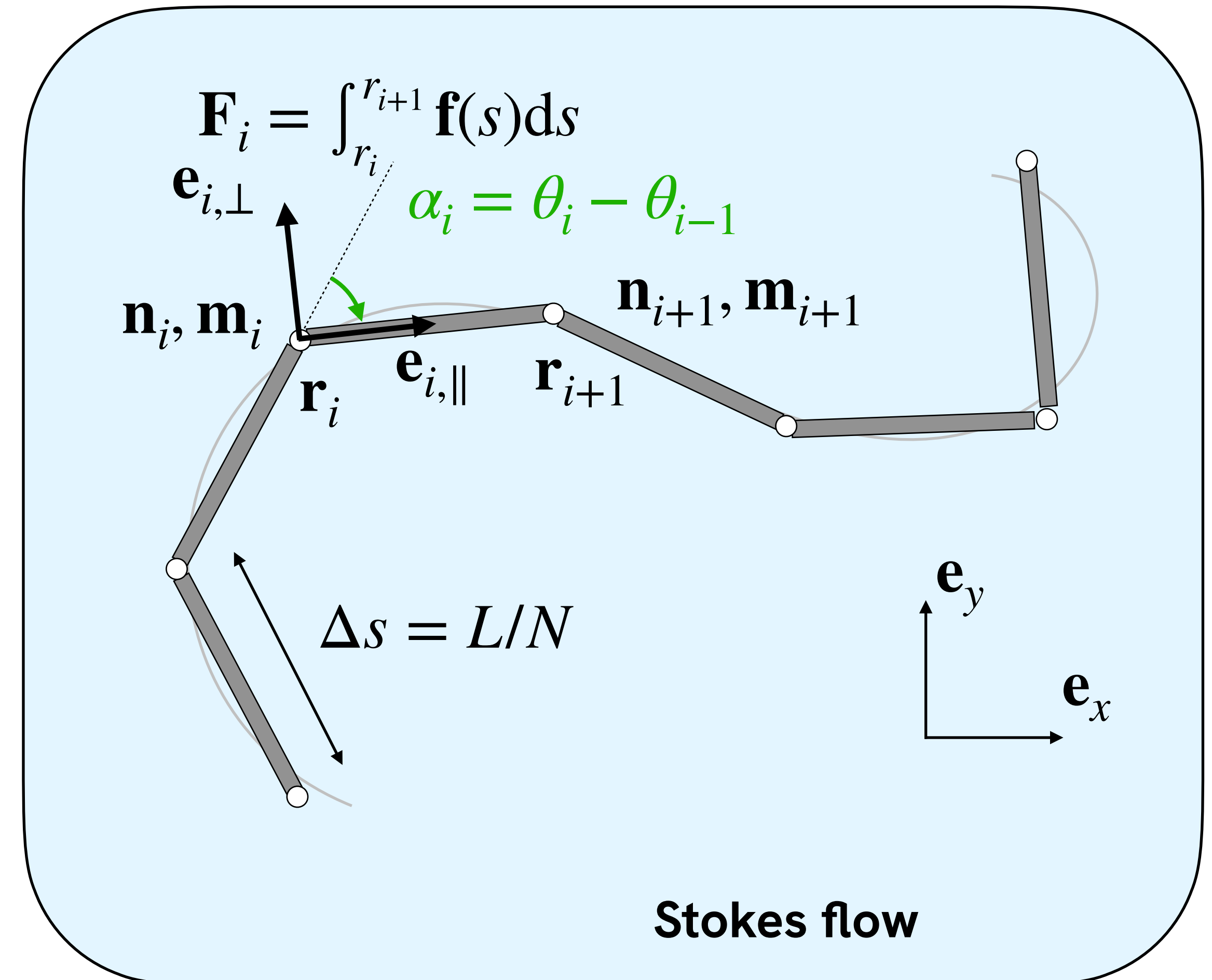
A “discrete” version: the N -link model



Elastohydrodynamics equations

A “discrete” version: the N -link model

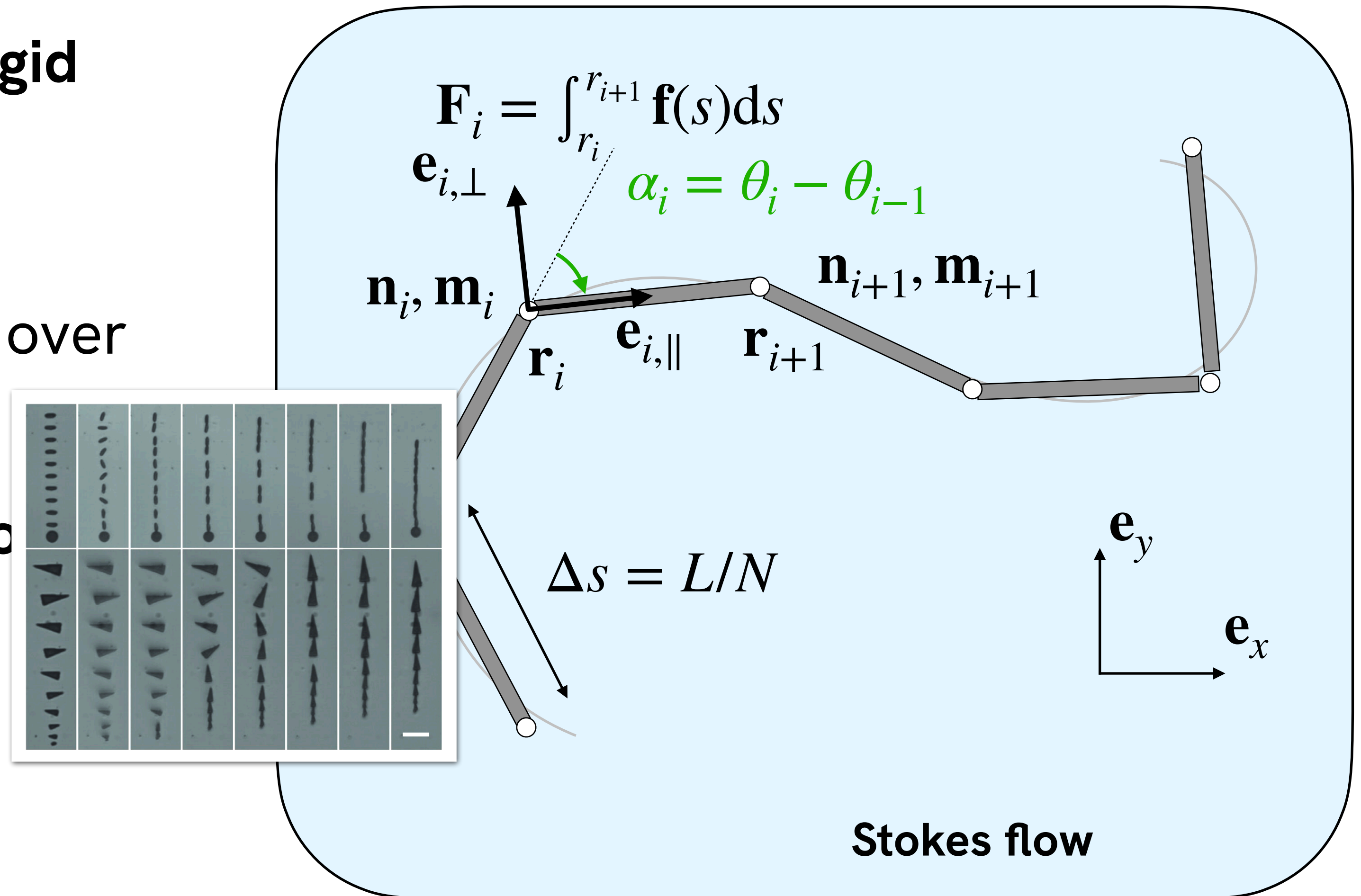
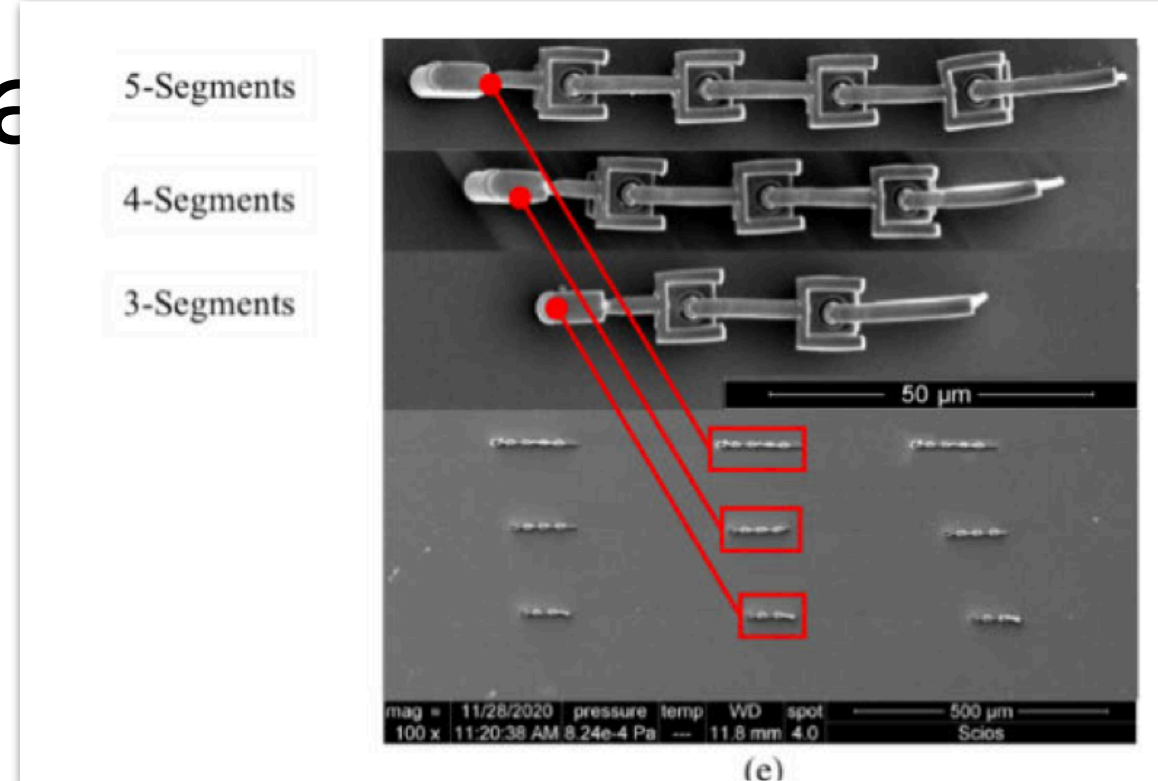
- Approximate the filament with **rigid segments** connected by **elastic junctions**
- Integrate the **drag force density** over each segment $\rightarrow$ ODE system
- A “mechanical” discretisation: **not** a discretization of arc length !



Elastohydrodynamics equations

A “discrete” version: the N -link model

- Approximate the filament with **rigid segments** connected by **elastic junctions**
- Integrate the **drag force density** over each segment $\rightarrow$ ODE system
- A “mechanical” discretisation: no discretization



Elastohydrodynamics equations

A “discrete” version: the N -link model

$$\forall i \in 1, \dots, N$$

(local force and torque balance)

$$\mathbf{n}_{i+1} - \mathbf{n}_i + \Delta s \mathbf{C}(\theta_i) \dot{\mathbf{r}}_{i+1/2} = 0$$

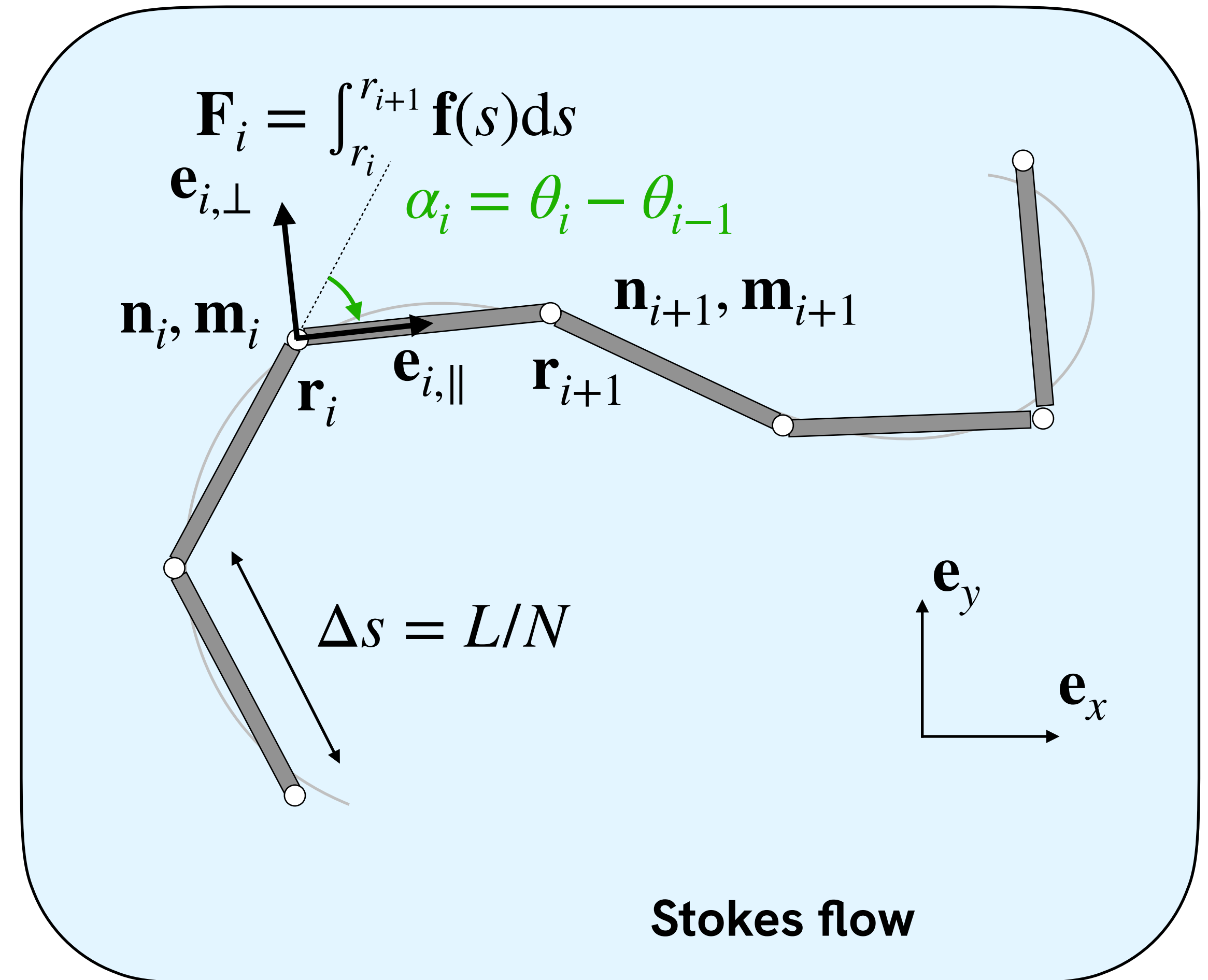
$$\mathbf{m}_{i+1} - \mathbf{m}_i + \Delta s \mathbf{e}_{i,\parallel} \times \frac{\mathbf{n}_{i+1} + \mathbf{n}_i}{2} - \frac{\Delta s^3}{12} c_{\perp} \dot{\theta}_i \mathbf{e}_z = 0$$

(Disc)

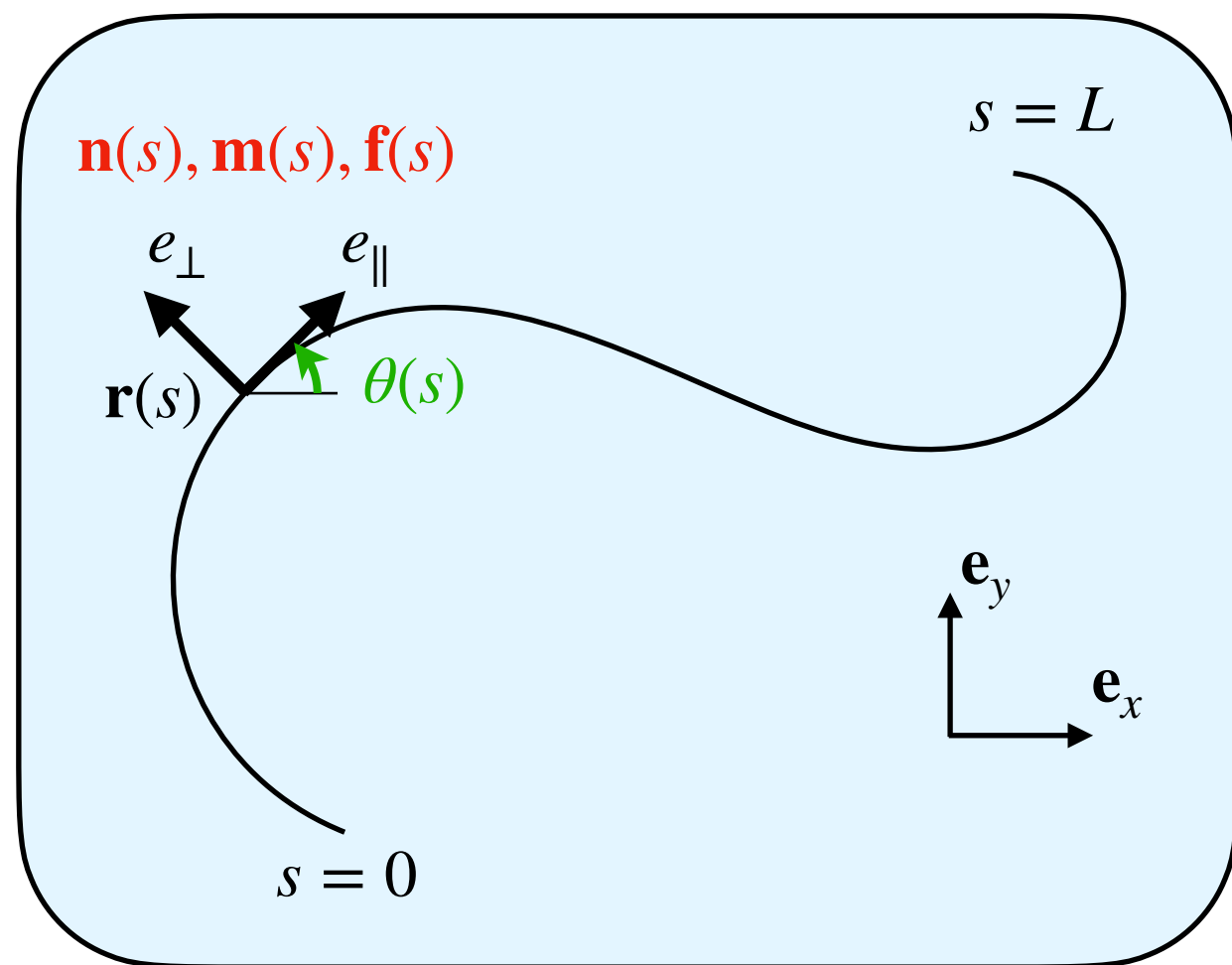
$$\mathbf{m}_{i+1} = \frac{E}{\Delta s} (\theta_{i+1} - \theta_i) \mathbf{e}_z \quad \text{(elastic constitutive law)}$$

$$|\mathbf{r}_{i+1} - \mathbf{r}_i| = \Delta s \quad \text{(inextensibility)}$$

+ B.C.



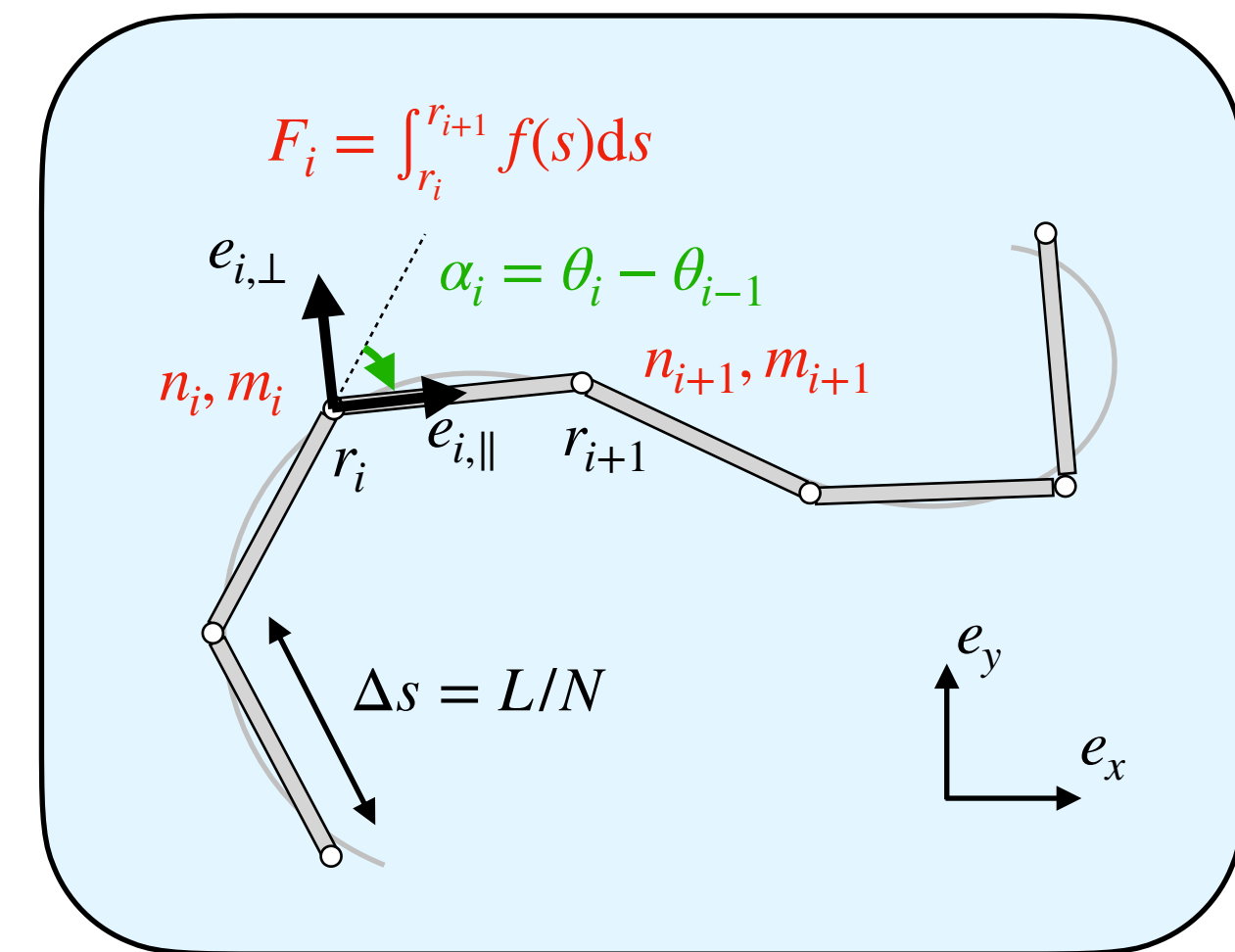
- Finite-dimensional, more efficient and robust numerically



Comparing both models

(Cont)

(Disc)



Local force balance

$$\partial_s \mathbf{n} + \mathbf{C}(\theta) \partial_t \mathbf{r} = 0$$

$$n_{i+1} - n_i + \Delta s \mathbf{C}(\theta_i) \dot{r}_{i+1/2} = 0$$

Local torque balance

$$\partial_s \mathbf{m} + \mathbf{r}_s \times \mathbf{n} = 0$$

$$m_{i+1} - m_i + \Delta s e_{i,\parallel} \times \frac{n_{i+1} + n_i}{2} - \frac{\Delta s^3}{12} c_{\perp} \dot{\theta}_i e_z = 0$$

Elastic constitutive law

$$\mathbf{m} = E \theta_s \mathbf{e}_z$$

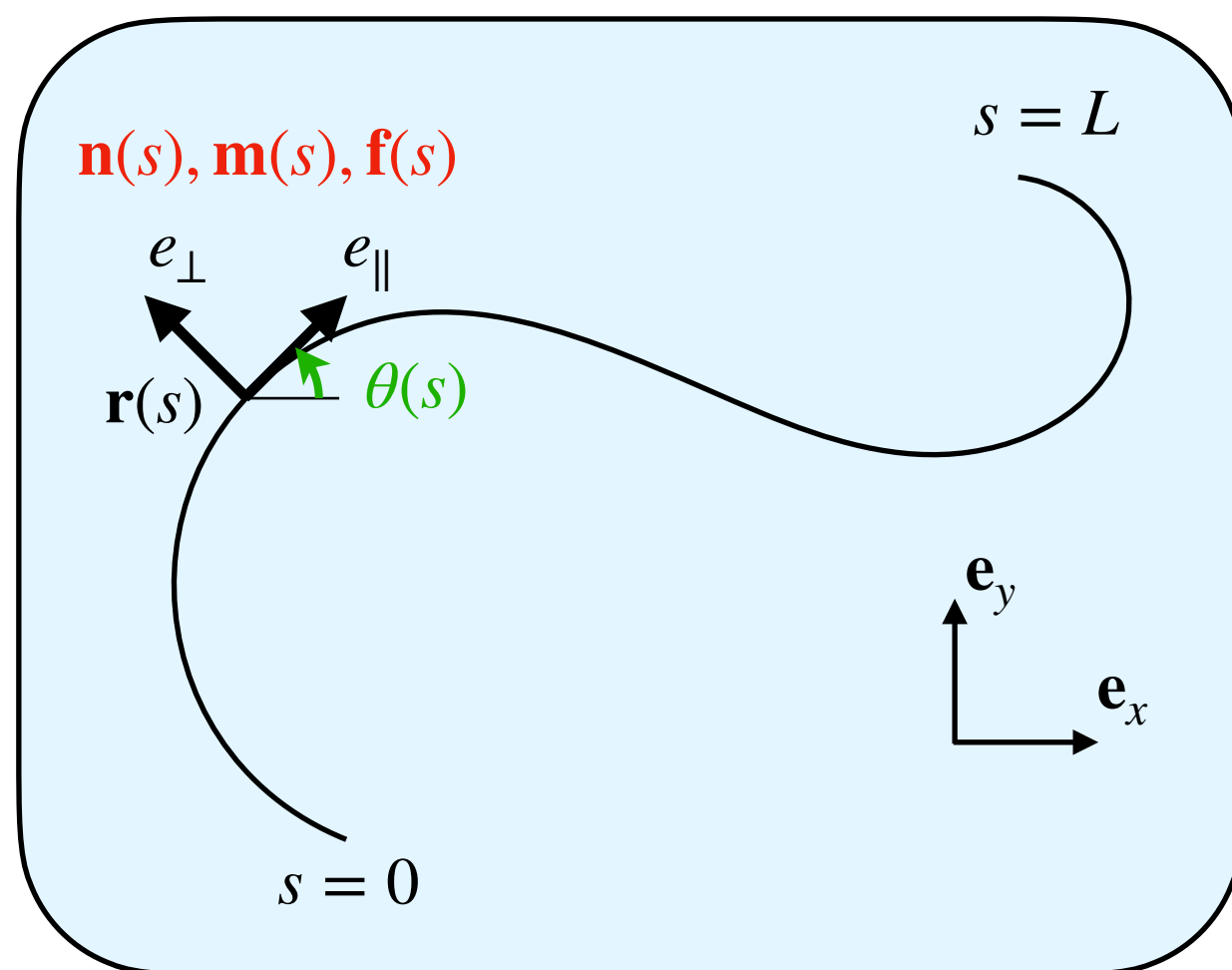
$$m_{i+1} = \frac{E}{\Delta s} (\theta_{i+1} - \theta_i) e_z$$

Inextensibility

$$|\mathbf{r}_s|^2 = 1$$

$$\frac{|r_{i+1} - r_i|}{\Delta s} = 1$$

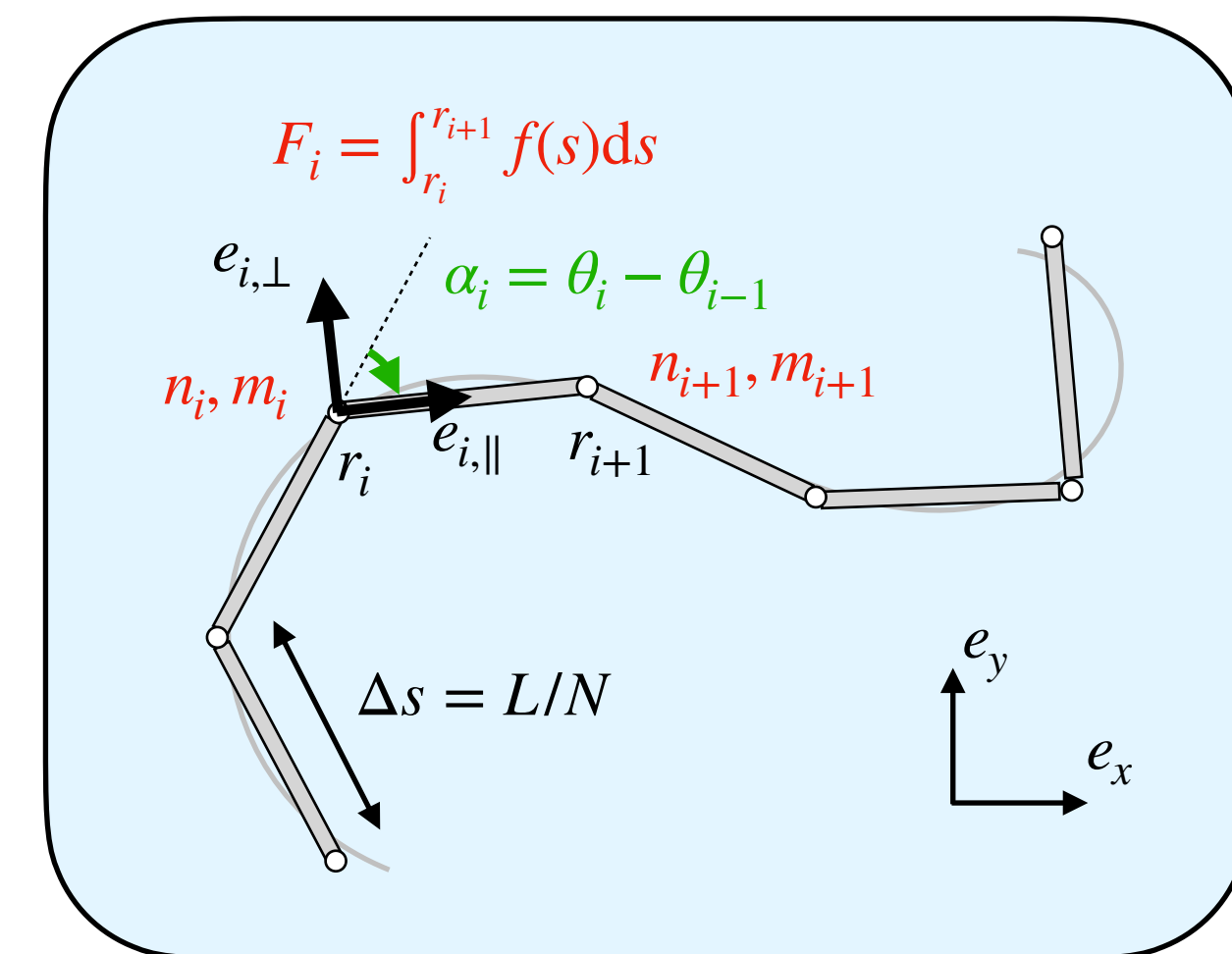
■ Convergence when $N \rightarrow \infty$?



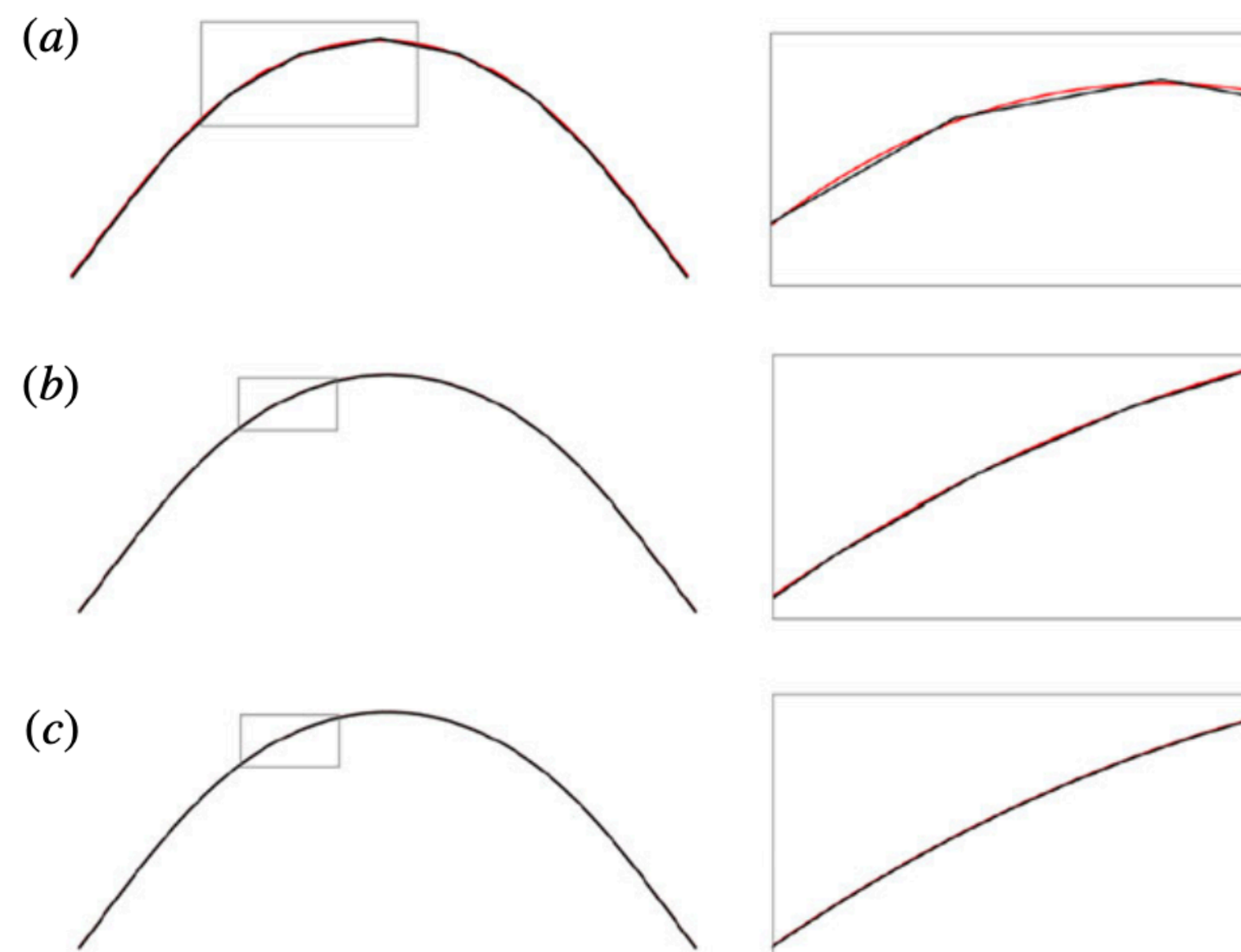
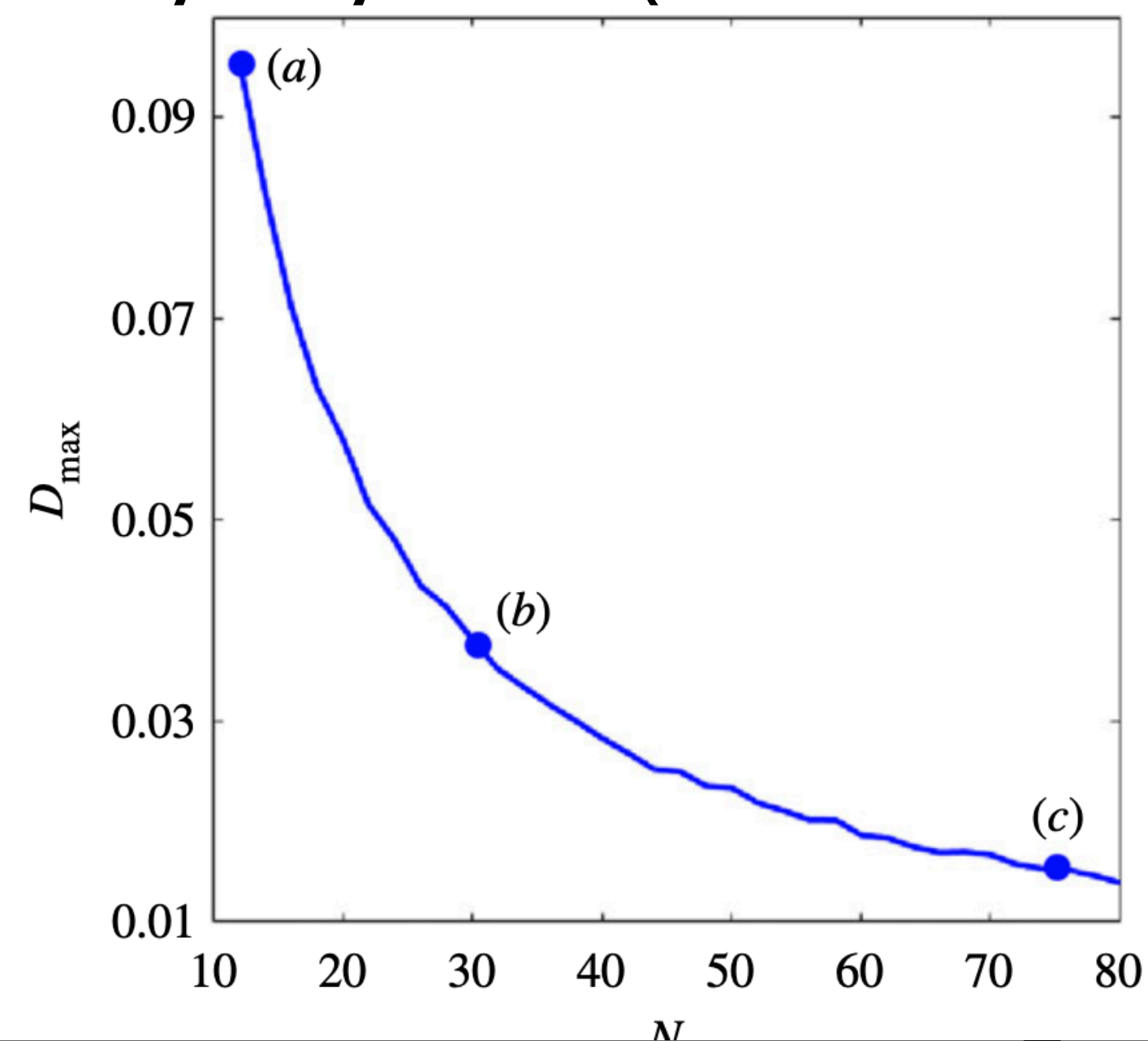
Comparing both models

(Cont)

(Disc)



numerically... easy to check (Moreau et al. J. R. Soc. Interface 15, 2018)



■ Convergence when $N \rightarrow \infty$?

$$\mathbf{C}(\theta_i) \dot{r}_{i+1/2} = 0$$

$$\frac{n_{i+1} + n_i}{2} - \frac{\Delta s^3}{12} c_{\perp} \dot{\theta}_i e_z = 0$$

$$(n_{i+1} - n_i) e_z$$

$$= 1$$

A convergence result

Theorem 1 (Alouges, Lefebvre-Lepot, Levillain, M., arXiv:2502.09988)

$$\text{(Disc)} \xrightarrow[N \rightarrow \infty]{} \text{(Cont)}$$

A convergence result

Theorem 1

(Cont)

I.C. $\mathbf{r}^0$ in $C^2(0,L)$

A convergence result

Theorem 1

(Disc): $\left((\mathbf{r}_i)_i, (\mathbf{n}_i)_i, (\mathbf{m}_i)_i \right)$

I.C. $(\mathbf{r}_i^0)_i$

←
define

(Cont)

I.C. $\mathbf{r}^0$ in $C^2(0,L)$

A convergence result

Interpolates

$$\mathbf{r}^h(t, s) = \sum_{i=1}^{N+1} \mathbf{r}_i(t) \phi_i(s), \quad m^h(t, s) = \sum_{i=1}^{N+1} m_i(t) \phi_i(s) \quad \mathbf{n}^h(t, s) = \sum_{i=1}^{N+1} \mathbf{n}_i(t) \phi_i(s)$$

Theorem 1

(Disc): $((\mathbf{r}_i)_i, (\mathbf{n}_i)_i, (\mathbf{m}_i)_i)$
I.C. $(\mathbf{r}_i^0)_i$

(Cont)
I.C. $\mathbf{r}^0$ in $C^2(0, L)$

A convergence result

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Theorem 1

(Disc): $((\mathbf{r}_i)_i, (\mathbf{n}_i)_i, (\mathbf{m}_i)_i)$
I.C. $(\mathbf{r}_i^0)_i$

(Cont)
I.C. $\mathbf{r}^0$ in $C^2(0, L)$

Assumptions (can always be satisfied)

- $\mathbf{r}^h(0, \cdot) \xrightarrow{N \rightarrow \infty} \mathbf{r}^0$ in $L^2(0, L)$
- $\exists C^0 > 0, \frac{E}{2} \sum_{i=1}^{N-1} h \left(\frac{\theta_{i+1}(0) - \theta_i(0)}{h} \right)^2 = C_h^0 \leq C^0$

A convergence result

Interpolates

$$\mathbf{r}^h(t, s) = \sum_{i=1}^{N+1} \mathbf{r}_i(t) \phi_i(s), \quad m^h(t, s) = \sum_{i=1}^{N+1} m_i(t) \phi_i(s) \quad \mathbf{n}^h(t, s) = \sum_{i=1}^{N+1} \mathbf{n}_i(t) \phi_i(s)$$

Theorem 1

(Disc): $((\mathbf{r}_i)_i, (\mathbf{n}_i)_i, (\mathbf{m}_i)_i)$

I.C. $(\mathbf{r}_i^0)_i$

I.C. $\mathbf{r}^0$

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Then,

$$\mathbf{r}^h \rightharpoonup \mathbf{r} \quad \text{in} \quad H^1(Q_T)$$

$$m^h \rightharpoonup m \quad \text{in} \quad L^2(0, T; H^1(0, L))$$

$$\mathbf{n}^h \rightharpoonup \mathbf{n} \quad \text{in} \quad L^2((0, T); H^1(0, L))$$

A convergence result

Interpolates

$$\mathbf{r}^h(t, s) = \sum_{i=1}^{N+1} \mathbf{r}_i(t) \phi_i(s), \quad m^h(t, s) = \sum_{i=1}^{N+1} m_i(t) \phi_i(s) \quad \mathbf{n}^h(t, s) = \sum_{i=1}^{N+1} \mathbf{n}_i(t) \phi_i(s)$$

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(Disc): $((\mathbf{r}_i)_i, (\mathbf{n}_i)_i, (\mathbf{m}_i)_i)$

I.C. $(\mathbf{r}_i^0)_i$

I.C. $\mathbf{r}^0$

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Then,

$$\mathbf{r}^h \rightharpoonup \mathbf{r} \quad \text{in } H^1(Q_T)$$

$$m^h \rightharpoonup m \quad \text{in } L^2(0, T; H^1(0, L))$$

$$\mathbf{n}^h \rightharpoonup \mathbf{n} \quad \text{in } L^2((0, T); H^1(0, L))$$

and $(\mathbf{r}, m, \mathbf{n})$ is weak sol. of (Cont)

Proof scheme

1. Show that $(\mathbf{r}^h, m^h, \mathbf{n}^h)$ are **bounded uniformly** in h in the relevant spaces

2. Extract a **weakly convergent subsequence** $(\mathbf{r}, m, \mathbf{n})$

 $\mathbf{r}^h$ m^h $\mathbf{n}^h$ $H^1(Q_T)$ $L^2(0,T; H^1(0,L))$ $L^2(0,T; H^1(0,L))$

Proof scheme

1. Show that $(\mathbf{r}^h, m^h, \mathbf{n}^h)$ are **bounded uniformly** in h in the relevant spaces

2. Extract a **weakly convergent subsequence** $(\mathbf{r}, m, \mathbf{n})$

3. Show that the limits are *(weakly)* solution to (Cont)

$\mathbf{r}^h$	$H^1(Q_T)$
m^h	$L^2(0,T; H^1(0,L))$
$\mathbf{n}^h$	$L^2(0,T; H^1(0,L))$

tricky! we need **additional bounds**, and **strong** convergence for some of them

Proof scheme

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3. Show that the limits are (weakly) solution to (Cont)

$\mathbf{r}^h$	$H^1(Q_T)$
m^h	$L^2(0,T; H^1(0,L))$
$\mathbf{n}^h$	$L^2(0,T; H^1(0,L))$
$\mathbf{r}_s^h$	$L^2(0,T; BV(0,L))$
$\dot{\mathbf{r}}_s^h$	$L^2(0,T; H^{-1}(0,L))$
$h\dot{\bar{\theta}}^h$	$L^2(Q_T)$
θ_s^h	$L^2(Q_T)$

tricky! we need **additional bounds**, and **strong** convergence for some of them

Detail on point 3

for all test functions φ in $C_c^\infty(0,L)$, ψ in $C_c^\infty(0,T)$

$$\int_0^T \int_0^L \left(\mathbf{C}(r_s(t,s)) \dot{r}(t,s) + n_s(t,s) \right) \varphi(s) \psi(t) \, ds \, dt = 0$$

force balance

$$\int_0^T \int_0^L \left(m_s(t,s) + r_s(t,s) \times n(t,s) \right) \varphi(s) \psi(t) \, ds \, dt = 0$$

moment balance

$$\int_0^T \int_0^L m^z(t,s) \varphi(s) \psi(t) \, ds \, dt = E \int_0^T \int_0^L (r_{ss}(t,s) \times r_s(t,s)) \cdot e_z \varphi(s) \psi(t) \, ds \, dt$$

constitutive equation

Detail on point 3

$$\int_0^T \int_0^L \left(\mathbf{C}(r_s(t, s)) \dot{\mathbf{r}}(t, s) + n_s(t, s) \right) \varphi(s) \psi(t) \, ds \, dt = 0$$

RFT:

$$\mathbf{C}(\theta(t, s)) \dot{\mathbf{r}} = -c_{\parallel}(\mathbf{e}_{\parallel} \cdot \dot{\mathbf{r}}) \mathbf{e}_{\parallel} - c_{\perp}(\mathbf{e}_{\perp} \cdot \dot{\mathbf{r}}) \mathbf{e}_{\perp}$$

- morally, product of $\mathbf{r}_s^2$ and $\dot{\mathbf{r}}$
- for weak convergence of a product $f \times g$, we need **strong** convergence of f and **weak** convergence of g
- so, we need **strong** convergence of $\mathbf{r}_s^2$ in $L^2 \rightarrow$ **strong** convergence of $\mathbf{r}_s$ in L^4

A BV estimate

$$\int_0^T \int_0^L \left(\mathbf{C}(r_s(t, s)) \dot{\mathbf{r}}(t, s) + n_s(t, s) \right) \varphi(s) \psi(t) \, ds \, dt = 0$$

RFT:

$$\mathbf{C}(\theta(t, s)) \dot{\mathbf{r}} = -c_{\parallel}(\mathbf{e}_{\parallel} \cdot \dot{\mathbf{r}}) \mathbf{e}_{\parallel} - c_{\perp}(\mathbf{e}_{\perp} \cdot \dot{\mathbf{r}}) \mathbf{e}_{\perp}$$

Theorem (Aubin 1963)

- Let X_0 , X and X_1 be three Banach spaces with $X_0 \subseteq X \subseteq X_1$.
- Suppose that $X_0 \Subset X$ and $X \hookrightarrow X_1$.
- Let $1 \leq p, q < +\infty$, let $W = \{u \in L^p(0, T; X_0) \mid \dot{u} \in L^q(0, T; X_1)\}$.
- Then the embedding of W into $L^p(0, T; X)$ is compact.

A BV estimate

$$\int_0^T \int_0^L \left(\mathbf{C}(r_s(t, s)) \dot{r}(t, s) + n_s(t, s) \right) \varphi(s) \psi(t) \, ds \, dt = 0$$

RFT:

$$\mathbf{C}(\theta(t, s)) \dot{\mathbf{r}} = -c_{\parallel}(\mathbf{e}_{\parallel} \cdot \dot{\mathbf{r}}) \mathbf{e}_{\parallel} - c_{\perp}(\mathbf{e}_{\perp} \cdot \dot{\mathbf{r}}) \mathbf{e}_{\perp}$$

Theorem (Aubin 1963)

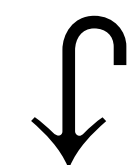
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- Then the embedding of W into $L^p(0, T; X)$ is compact.

$$H^{-1}(0, L)$$



compact
embedding

$$L^4(0, L)$$



continuous
embedding

$$BV(0, L)$$

Key formula: energy dissipation for (Disc)

Proposition:

$$\frac{1}{2} \frac{d}{dt} \left[E \sum_{i=1}^{N-1} \Delta s \left(\frac{\theta_{i+1} - \theta_i}{\Delta s} \right)^2 \right] + \sum_{i=1}^N \frac{\Delta s^3}{12} c_{\perp} \dot{\theta}_i^2 - \sum_{i=1}^N \Delta s \mathbf{C}(\theta_i) \dot{\mathbf{r}}_{i+\frac{1}{2}} \cdot \dot{\mathbf{r}}_{i+\frac{1}{2}} = 0.$$

Elastic energy

Hydrodynamic drag
dissipation

Bounds

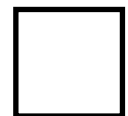
$\mathbf{r}^h$	$H^1(Q_T)$	inextensibility + energy dissipation + Poincaré-Wirtinger
m^h	$L^2(0,T; H^1(0,L))$	energy dissipation
$\mathbf{n}^h$	$L^2(0,T; H^1(0,L))$	energy dissipation
$\mathbf{r}_s^h$	$L^2(0,T; BV(0,L))$	compute total variation + assumption on bending energy
$\dot{\mathbf{r}}_s^h$	$L^2(0,T; H^{-1}(0,L))$	consequence of (1)
$h\dot{\theta}^h$	$L^2(Q_T)$	energy dissipation
θ_s^h	$L^2(Q_T)$	energy dissipation

Convergences

$\mathbf{r}^h$	$\mathbf{r}$	$H^1(Q_T)$	weak	immediate from boundedness
$\mathbf{r}^h$	$\mathbf{r}$	$L^2(Q_T)$	strong	Rellich-Kondrachov theorem
$\mathbf{r}_s^h$	$\mathbf{r}_s$	$L^2(0,T; L^p(0,L))$	strong	compact embedding of BV + Aubin-Lions-Simon theorem
$\hat{\mathbf{r}}^h$	$\mathbf{r}$	$L^2(0,T; H^1(0,L))$	strong	connect to bounds on m and θ
$\hat{\mathbf{r}}^h$	$\mathbf{r}$	$L^2(0,T; H^2(0,L))$	weak	
$\mathbf{n}^h$	$\mathbf{n}$	$L^2(0,T; H^1(0,L))$	weak	immediate from boundedness
$\mathbf{m}^h$	$\mathbf{m}$	$L^2(0,T; H^1(0,L))$	weak	immediate from boundedness
$\bar{\mathbf{m}}^h$	$\mathbf{m}$	$L^2(0,T; L^p(0,L))$	weak	show that $\ \mathbf{m} - \bar{\mathbf{m}}\ _{L^2(Q_T)} \rightarrow 0$
$h^2 \dot{\bar{\theta}}^h$	0	$L^2(Q_T)$	strong	immediate from boundedness
θ_s^h	α	$L^2(Q_T)$	weak	immediate from boundedness

Convergences

$\mathbf{r}^h$	$\mathbf{r}$	$H^1(Q_T)$	weak	immediate from boundedness
$\mathbf{r}^h$	$\mathbf{r}$	$L^2(Q_T)$	strong	Rellich-Kondrachov theorem
$\mathbf{r}_s^h$	$\mathbf{r}_s$	$L^2(0,T; L^p(0,L))$	strong	compact embedding of BV + Aubin-Lions-Simon theorem
$\hat{\mathbf{r}}^h$	$\mathbf{r}$	$L^2(0,T; H^1(0,L))$	strong	connect to bounds on m and θ
$\hat{\mathbf{r}}^h$	$\mathbf{r}$	$L^2(0,T; H^2(0,L))$	weak	
$\mathbf{n}^h$	$\mathbf{n}$	$L^2(0,T; H^1(0,L))$	weak	immediate from boundedness
$\mathbf{m}^h$	$\mathbf{m}$	$L^2(0,T; H^1(0,L))$	weak	immediate from boundedness
$\bar{\mathbf{m}}^h$	$\mathbf{m}$	$L^2(0,T; L^p(0,L))$	weak	show that $\ \mathbf{m} - \bar{\mathbf{m}}\ _{L^2(Q_T)} \rightarrow 0$
$h^2 \dot{\bar{\theta}}^h$	0	$L^2(Q_T)$	strong	immediate from boundedness
θ_s^h	α	$L^2(Q_T)$	weak	immediate from boundedness



Back to the theorem

Interpolates

$$\mathbf{r}^h(t, s) = \sum_{i=1}^{N+1} \mathbf{r}_i(t) \phi_i(s), \quad m^h(t, s) = \sum_{i=1}^{N+1} m_i(t) \phi_i(s) \quad \mathbf{n}^h(t, s) = \sum_{i=1}^{N+1} \mathbf{n}_i(t) \phi_i(s)$$

Theorem 1 (Alouges, Lefebvre-Lepot, Levillain, M., arXiv:2502.09988)

(Disc): $((\mathbf{r}_i)_i, (\mathbf{n}_i)_i, (\mathbf{m}_i)_i)$

I.C. $\mathbf{r}^0$

I.C. $(\mathbf{r}_i^0)_i$

Assumptions (can always be satisfied)

- $\mathbf{r}^h(0, \cdot) \xrightarrow{N \rightarrow \infty} \mathbf{r}^0$ in $L^2(0, L)$
- $\exists C^0 > 0, \frac{E}{2} \sum_{i=1}^{N-1} h \left(\frac{\theta_{i+1}(0) - \theta_i(0)}{h} \right)^2 = C_h^0 \leq C^0$

Then,

$$\mathbf{r}^h \rightharpoonup \mathbf{r} \quad \text{in} \quad H^1(Q_T)$$

$$m^h \rightharpoonup m \quad \text{in} \quad L^2(0, T; H^1(0, L))$$

$$\mathbf{n}^h \rightharpoonup \mathbf{n} \quad \text{in} \quad L^2((0, T); H^1(0, L))$$

and $(\mathbf{r}, m, \mathbf{n})$ is sol. of (Cont)

Existence?

Interpolates

$$\mathbf{r}^h(t, s) = \sum_{i=1}^{N+1} \mathbf{r}_i(t) \phi_i(s), \quad m^h(t, s) = \sum_{i=1}^{N+1} m_i(t) \phi_i(s) \quad \mathbf{n}^h(t, s) = \sum_{i=1}^{N+1} \mathbf{n}_i(t) \phi_i(s)$$

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$$m^h \rightharpoonup m \quad \text{in} \quad L^2(0, T; H^1(0, L))$$

$$\mathbf{n}^h \rightharpoonup \mathbf{n} \quad \text{in} \quad L^2((0, T); H^1(0, L))$$

and $(\mathbf{r}, m, \mathbf{n})$ is sol. of (Cont)

Existence of solutions to (Disc)?

(Disc): $\left((\mathbf{r}_i)_i, (\mathbf{n}_i)_i, (\mathbf{m}_i)_i \right)$

I.C. $(\mathbf{r}_i^0)_i$

Theorem 2 (*Well-posedness of (Disc)*)

For all initial data, there exists a unique global solution of (Disc) in $C^1(\mathbb{R}_+)$.

Existence of solutions to (Disc)?

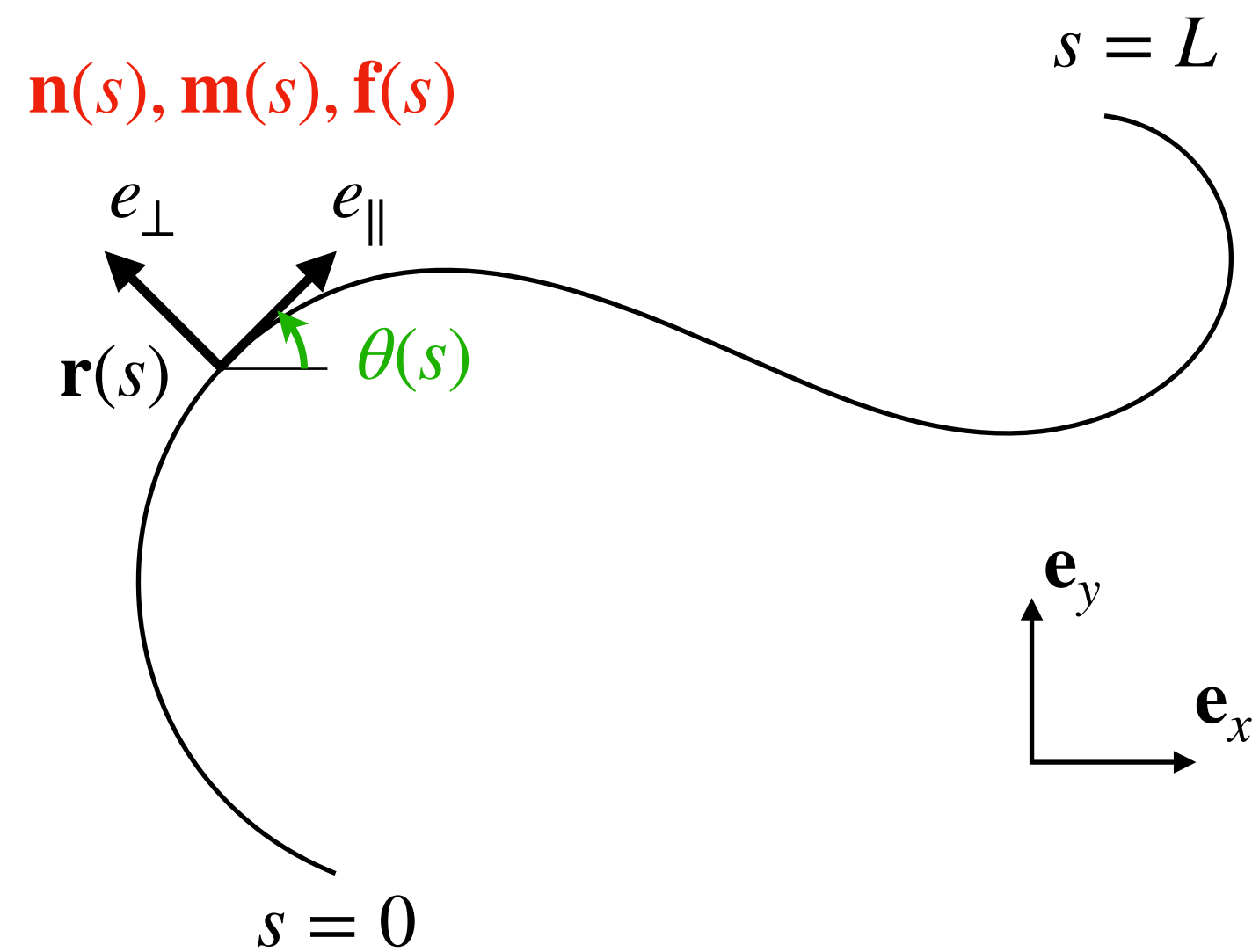
(Disc): $((\mathbf{r}_i)_i, (\mathbf{n}_i)_i, (\mathbf{m}_i)_i)$

I.C. $(\mathbf{r}_i^0)_i$

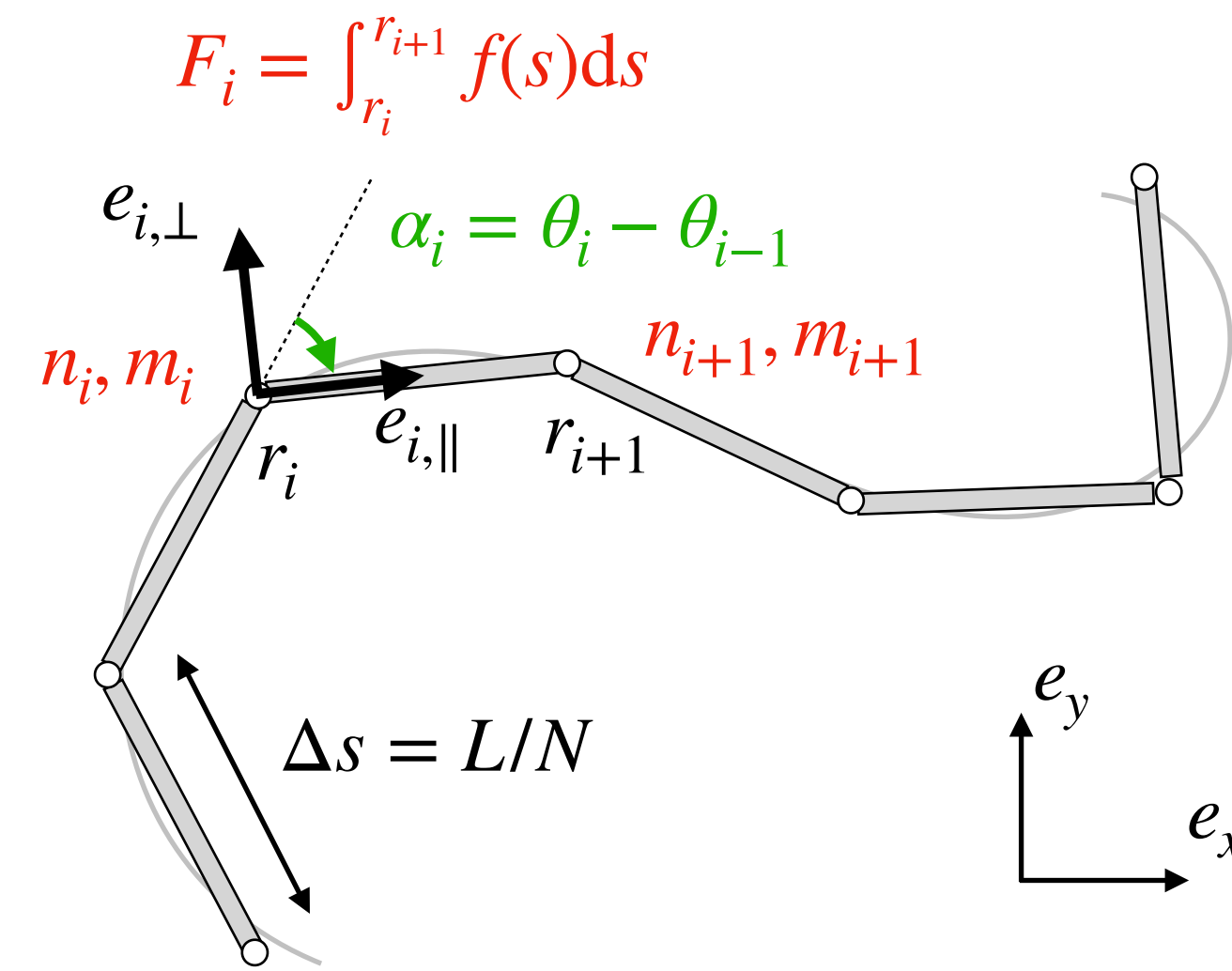
Theorem 2 (*Well-posedness of (Disc)*)

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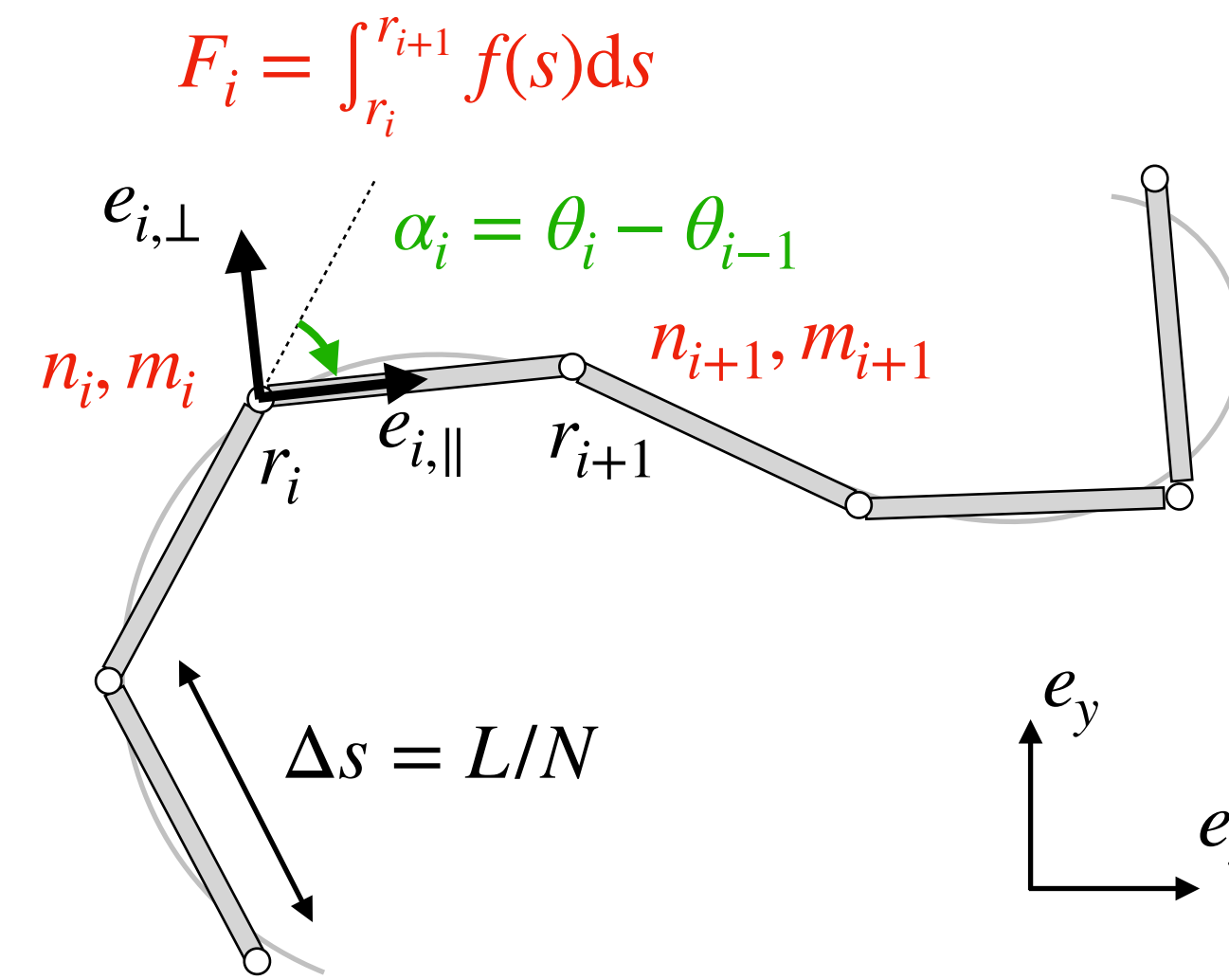
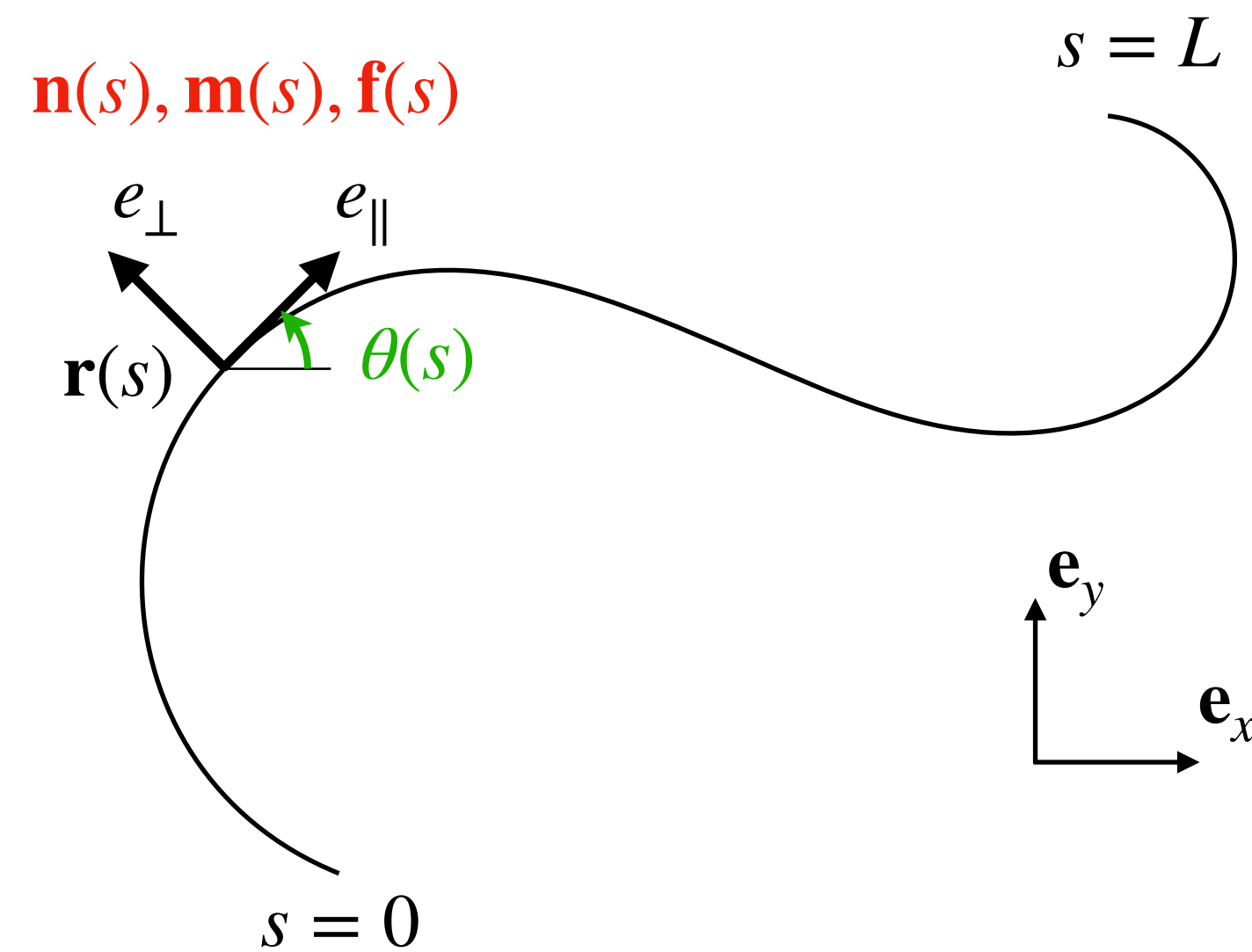
- Differential algebraic system: check invertibility + bounds to have global existence through Cauchy-Lipschitz theorem



Theorem 1: Weak convergence of (Disc) towards (Cont)



Theorem 2: Well-posedness of the N -link model



Theorem 1: Weak convergence of (Disc) towards (Cont)

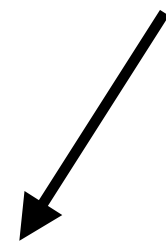
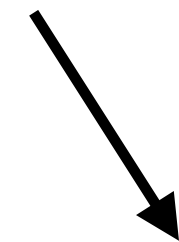
Theorem 2: Well-posedness of the N -link model

Well-posedness of (Cont)?

Well-posedness of (Cont)?

Theorem 1: Weak convergence of (Disc) towards (Cont)

Theorem 2: Existence and uniqueness of solutions for the N -link model



Corollary: Existence of solutions for the continuous model

but... up to the extraction of subsequences! $\rightarrow$ uniqueness is not given for free

Well-posedness of (Cont)?

Theorem 1: Weak convergence of (Disc) towards (Cont)

Theorem 2: Existence and uniqueness of solutions for the N -link model



Corollary: Existence of solutions for the continuous model in $H^1([0,T]; H^1) \cap L^2([0,T], H^2)$

Theorem: (Mori, Ohm 2023) Existence and uniqueness of solutions *globally in time* for *small initial data* in $C^0([0,T]; L^2) \cap C^0((0,T]; H^1)$

Well-posedness of (Cont)?

Theorem 1: Weak convergence of (Disc) towards (Cont)

Theorem 2: Existence and uniqueness of solutions for the N -link model



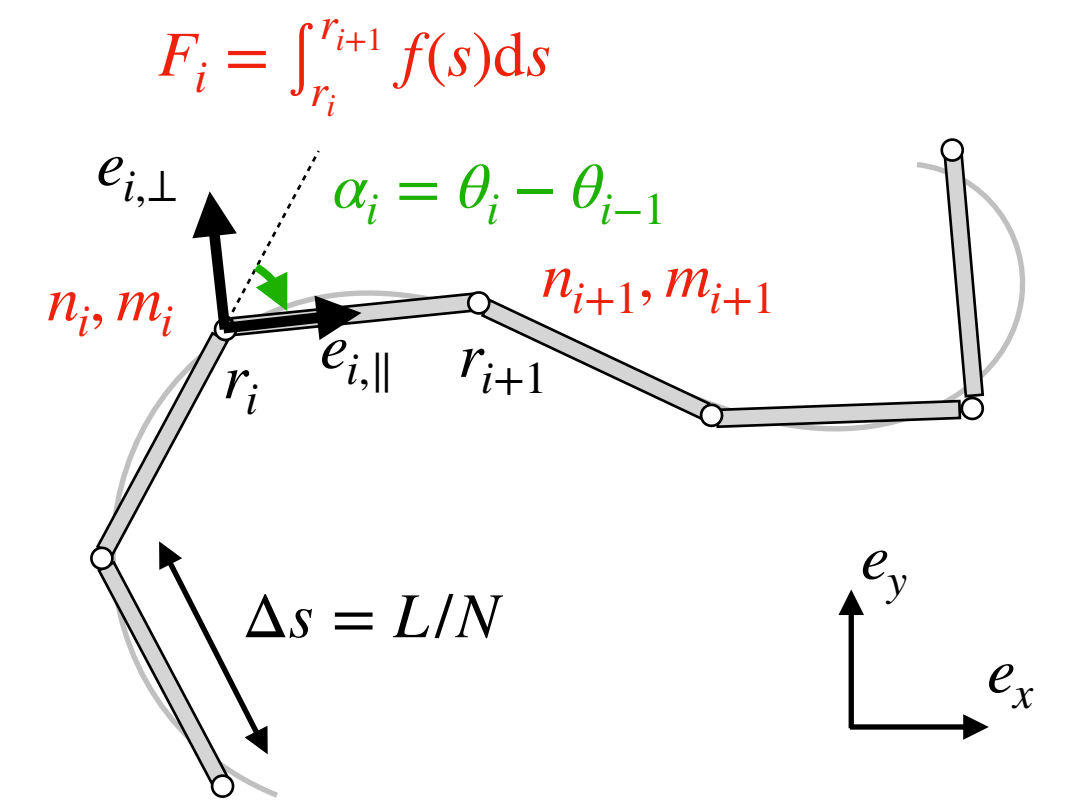
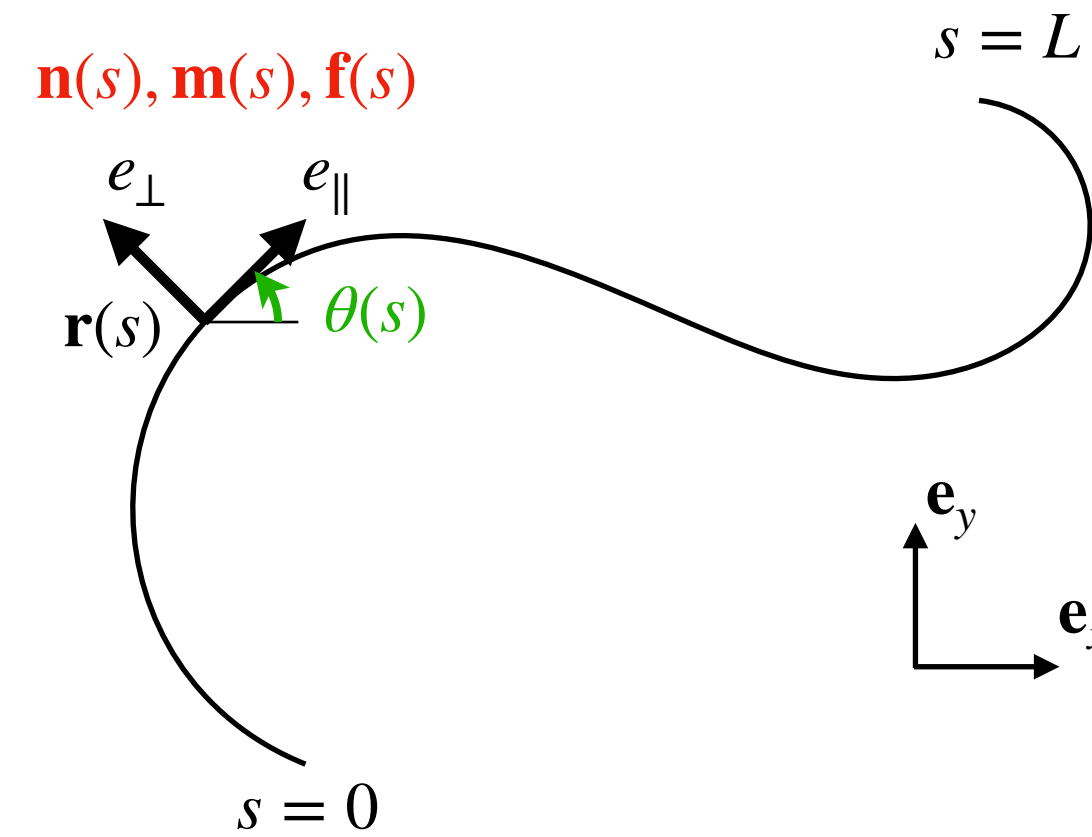
Corollary: Existence of solutions for the continuous model in $H^1([0,T]; H^1) \cap L^2([0,T], H^2)$

Theorem: (Mori, Ohm 2023) Existence and uniqueness of solutions *globally in time* for *small initial data* in $C^0([0,T]; L^2) \cap C^0((0,T]; H^1)$

Theorem: (Albritton, Ohm 2025) Global existence and uniqueness of solutions in $L^2([0,T]; H^2)$

Summary and outlook

arXiv:2502.09988



Theorem 1: Weak convergence of (Disc) towards (Cont)

Theorem 2: Well-posedness of the N -link model

Corollary: Existence of solutions for the continuous model

Extensions

- Internal activity, curvature at rest
- Refined hydrodynamics
- 3D, non-Newtonian fluids, multiple filaments...



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